

# 1        **Ionospheric $D$ region: Characteristics near Dawn and Dusk**

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## 8        **Key Points:**

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- ‘Wait’  $D$  region electron number density parameters, height  $H'$  and sharpness

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$\beta$ , are determined through dawn and dusk for the first time

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- At mid-latitudes  $H'$  shows a clear minimum at dawn consistent with the

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release of cosmic ray generated electrons accumulated overnight

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- $H'$  and  $\beta$  through dawn and dusk are a delicate balance between  $D$ -region

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ionization generated by solar EUV and galactic cosmic rays

## 15 **Abstract**

16 The characteristics of very low frequency (VLF) radio wave propagation in the Earth-  
17 ionosphere waveguide are determined particularly through dawn and dusk using  
18 phase and amplitude measurements of man-made signals propagating below the  
19 ionospheric *D* region. For the first time variations of ‘Wait’ height and sharpness  
20 parameters,  $H'$  and  $\beta$ , have been determined for dawn and dusk conditions. These  
21 measurements provide observational data to constrain *D* region modeling efforts,  
22 extending the capabilities of VLF propagation monitoring for geophysical phenomena  
23 such as lightning, solar flares, and energetic particle precipitation. At mid-latitudes,  $H'$   
24 varied from  $\sim 85$  km at night, then, starting from solar zenith angle (SZA)  $\sim 97.5^\circ$ ,  
25 rapidly down to  $\sim 73$  km at dawn (SZA  $\sim 90^\circ$ ), then back up to  $\sim 78$  km at SZA  $\sim 75^\circ$   
26 and then down to the appropriate noon value for the latitude (and season). In contrast,  
27 from noon through dusk to night,  $H'$  varied essentially monotonically from  $\sim 70$ - $75$  km  
28 through  $\sim 80$  km to  $\sim 85$  km. At low latitudes no dawn minimum in  $H'$  was observed,  
29 due to the reduced effect of galactic cosmic rays. Sharpness,  $\beta$ , varied from its  
30 nighttime value of  $\sim 0.6$  km $^{-1}$  down to a minimum of  $\sim 0.25$  km $^{-1}$  at SZA  $\sim 85^\circ$  near  
31 dusk or  $\sim 75^\circ$  near dawn, rising again to (SZA-dependent) noon values of  $\sim 0.35$ - $0.5$   
32 km $^{-1}$ . The results are interpreted through the geophysical effects controlling *D* region  
33 electrons, including the daytime dominant role of solar Lyman- $\alpha$  from low to mid-  
34 latitudes, and the greater role of galactic cosmic rays at increasingly higher mid-  
35 latitudes.

## 36 **1. Introduction**

37 Very Low Frequency (VLF) radio waves can propagate over very long distances  
38 (thousands of km) in the Earth-ionosphere waveguide bounded below by the Earth's  
39 surface (oceans/ground) and bounded above by the lowest edge of the Earth's  
40 ionosphere (the lower *D* region) at heights ~70 km by day and ~85 km by night. VLF  
41 radio propagation in the ionosphere is essentially controlled by free electrons; both  
42 positive and negative ions occur in comparable concentrations but have no significant  
43 effect at VLF because they are so much more massive than the free electrons (by  
44 factors of >10,000).

45 In the quiet *D* region there are two principal sources of free electrons. The main one at  
46 low and middle latitudes is UV from the daytime Sun, mainly Lyman- $\alpha$ , ionizing the  
47 minor neutral constituent NO (to  $\text{NO}^+ + e^-$ ); this is generally important (and  
48 dominant) only above 65-70 km altitude because Lyman- $\alpha$  is absorbed by neutral  $\text{O}_2$   
49 below these heights (Banks & Kockarts, 1973). This absorption by  $\text{O}_2$  also results in  
50 the ionization of NO being solar zenith angle (SZA) dependent and thus dependent on  
51 time of day and latitude.

52 The other important ionizing source is galactic cosmic rays (GCR) which ionize all  
53 the constituents in the neutral atmosphere, day and night, 24 hours a day. This  
54 ionizing process tends to be dominant below heights of 65-70 km at low to mid-  
55 latitudes because of the shielding effect of the Earth's geomagnetic field. In contrast,  
56 at mid- to high latitudes GCR ionization tends to dominate up to greater heights, 70-  
57 75 km or more, because of the lower shielding by the Earth's geomagnetic field there  
58 and the Sun being lower in the sky (i.e., higher SZA) resulting in more absorption of  
59 Lyman- $\alpha$  by  $\text{O}_2$ . At night virtually all of the free electrons below ~75 km become  
60 effectively removed by rapidly attaching to the copious neutral  $\text{O}_2$  molecules at these

61 low altitudes forming negative ions. When daylight returns, visible light from the Sun  
62 releases these electrons again from the negative ions (Peterson, 1976; Reid, 1987;  
63 Thomas & Bowman, 1986; Thomas & Harrison, 1970; Verronen et al. 2006; see also  
64 Banks & Kockarts, 1973).

65 VLF propagation in the Earth-Ionosphere waveguide has been found to be normally  
66 remarkably stable in unperturbed conditions. Phase and amplitude perturbations in the  
67 propagation have thus been able to be used extensively to monitor external  
68 perturbations (in energy and height) such as from energetic particle (particularly  
69 electron) precipitation from the Earth's radiation belts into the top of the waveguide  
70 (e.g., Rodger et al., 2010). For convenience in making VLF propagation calculations,  
71 VLF observations have been used successfully to characterize the electron densities in  
72 the lower  $D$  region by determining the 'Wait' (Wait & Spies, 1964) height and  
73 sharpness parameters,  $H'$  and  $\beta$ , under a variety of daytime and nighttime conditions.  
74 These have included daytime at low latitudes (Thomson, 2010; Thomson et al., 2014),  
75 daytime at mid-latitudes (Thomson et al., 2011a, 2017), and nighttime (Thomson et  
76 al., 2007; Thomson & McRae, 2009), all of which used narrow-band VLF  
77 transmissions from man-made transmitters. Broadband VLF from lightning has been  
78 used to make similar measurements mainly over the US continental land mass,  
79 (Cummer et al., 1998 and Cheng et al., 2005 at night; Han & Cummer, 2010 by day).

80 VLF propagation techniques have recently been used to determine  $D$  region  
81 parameters by Kumar & Kumar (2020) and Chand & Kumar (2021) using VLF man-  
82 made transmitters received in the South Pacific. VLF techniques have also recently  
83 been used by McCormick et al. (2021) testing a new four-parameter  $H'/\beta$  model for  
84 the  $D$  region, by Zhou et al. (2021) using timing of first and second lightning hops,  
85 and by Chowdhury et al. (2021) from the International Reference Ionosphere (IRI-

86 2016, Bilitza, 2017) with VLF validation. Additionally, VLF techniques have been  
87 used by Rozhnoi et al. (2021), Barman et al. (2024), and Basak et al. (2024) to study  
88 solar eclipses, by Macotela et al. (2021) reporting a daytime spring-fall amplitude  
89 asymmetry, and by Xu et al. (2021) and Worthington & Cohen (2021) using VLF  
90 trans-ionospheric propagation together with the rocket-based Faraday International  
91 Reference Ionosphere electron density model (FIRI, e.g., Friedrich et al., 2018).

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93 The determination and validation of energetic electron precipitation characteristics  
94 impacting the atmosphere has been identified as a key component of solar forcing  
95 descriptions in coupled-climate modeling studies (Matthes et al., 2017; Funke et al.,  
96 2024). The use of VLF propagation measurements to calculate electron precipitation  
97 fluxes has been an important factor in these efforts. Similarly effects of solar flares  
98 have been extensively monitored (e.g., Thomson & Clilverd, 2001; Thomson et al.,  
99 2005), as have occasional extra-galactic gamma-ray bursts (e.g., Fishman & Inan,  
100 1988; Pal et al., 2023). Both the World-Wide Lightning Location Network (WWLLN,  
101 e.g., Rodger et al., 2005) and the Global Lightning Dataset (GLD360, e.g., Said et al.,  
102 2010) rely on VLF radio waves radiated from lightning flashes propagating up to  
103 many thousands of km to the VLF receivers of their global networks. The world's  
104 great naval powers communicate with their submarines from some of the same large  
105 man-made VLF transmitters we use here. They make use of the large horizontal  
106 ranges (many thousands of km) and the ability of very low frequencies to penetrate  
107 seawater so allowing their submarines to remain submerged.

108 However, such efforts have been hampered by the difficulties in understanding  
109 propagation conditions near sunrise and sunset (Clilverd et al., 2010; Neal et al.,  
110 2015). A lack of detailed knowledge of VLF propagation conditions around sunrise

111 and sunset has also restricted electron precipitation validation studies (Clilverd et al.,  
112 2020). For many years, distinctive VLF anomalies have been observed on (long)  
113 propagation paths that travel across a dawn/dusk terminator, particularly at dawn  
114 (Crombie, 1964). These involve clear successive minima in amplitude, known as  
115 modal minima, likely caused by destructive interference between new modes  
116 generated by the *D* region characteristics rapidly changing with distance as the waves  
117 travel across the terminator. These amplitude minima are often accompanied by rapid  
118 phase changes known as cycle slips. Recently Chand & Kumar (2017) reported and  
119 discussed similar observations on long VLF paths recorded at Fiji in the South  
120 Pacific. Quantitative explanations of these anomalies have been hindered by lack of *D*  
121 region propagation characteristics at dawn and dusk. In many of these cases the VLF  
122 waves will at times travel from a nighttime region of the Earth to a daytime region or  
123 vice-versa and so will propagate through a dawn/dusk transition requiring a model of  
124 the appropriate dawn/dusk ionosphere in this transition region for predicting phase  
125 and amplitude (as a function of time) at the receiver, to compare with observations.  
126 However, very little attention has been given to characterize the *D* region in the dawn  
127 and dusk transition regions between day and night. Hence observations of the *D*  
128 region near dawn and dusk, as here, are very desirable for modeling and testing  
129 potential mechanisms in these transition regions.

130 A key aim of the current work is to provide the measurements to facilitate the  
131 modelling of dawn and dusk propagation conditions. Such measurements will extend  
132 the capabilities of VLF propagation monitoring for geophysical phenomena such as  
133 lightning, solar flares, solar eclipses, geomagnetic storms, substorms, and energetic  
134 particle precipitation. In order to accurately model propagation conditions over a wide  
135 range of solar zenith angle (SZA), observations from specific transmitter–receiver

136 paths with near-constant SZA along the whole path are analysed. Interpretation is  
137 made using long-wave propagation codes. The selection and modelling of these paths  
138 are described in Section 2. The delicate balance between solar UV and galactic  
139 cosmic ray effects is investigated within three different latitude bands. Section 3  
140 presents the results for mid-latitude paths, with detailed analysis of specific paths.  
141 Section 4 present results for the low-latitude paths, while Section 5 presents results  
142 for the high latitude paths. The results are discussed in Section 6, including  
143 comparisons between dawn/dusk conditions relative to the overall diurnal variations,  
144 and existing rocket measurements. Summary and conclusions are given in Section 7.  
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## 150 **2. Methods**

### 151 **2.1 Selecting paths for VLF Observations**

152 For a typical VLF path, at a fixed time of day, the SZA varies along the path from  
153 transmitter to receiver making it not straightforward to determine the characteristics at  
154 a particular SZA. Paths greater in length than ~1000 km typically needed to be  
155 specifically chosen such that, at some convenient time of year when suitable VLF  
156 data was available, the SZA varied very little along the path at fixed times. This is  
157 particularly the case for dawn or dusk, so that the characteristics of the path at  
158 particular SZA could be determined. This meant that the path needed to align at least  
159 approximately with the sunrise or sunset, day/night-terminator at that time of year.  
160 The ‘convenient time of year’ was preferably in the <~6 months of the year closest to  
161 summer solstice for calibrating the path’s VLF propagation using mid-day  
162 observations. The paths so chosen are shown in Figure 1. These include two very  
163 short paths (~300 km), from NPM in Hawaii (Thomson et al., 2014) and from NWC,  
164 North West Cape, Australia (Thomson et al., 2012 & 2014). For these the SZA varies  
165 very little along their short paths (which also happen to be at low latitudes). The  
166 receivers at Eskdalemuir, St. John’s, Reykjavik and Dunedin are part of the  
167 AARDDVARK network (Antarctic-Arctic Radiation-belt (Dynamic) Deposition-VLF  
168 Atmospheric Research Konsortium: e.g., Clilverd et al., 2009;  
169 [http://www.physics.otago.ac.nz/space/AARDDVARK\\_homepage.htm](http://www.physics.otago.ac.nz/space/AARDDVARK_homepage.htm) ).

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### 171 **2.2 Determining $D$ region parameters, $H'$ and $\beta$ , from the VLF Observations**

172 US Navy modal waveguide codes, ModeFinder (Morfitt & Shellman, 1976) and  
173 LWPC (Long Wavelength Propagation Capability, Ferguson & Snyder, 1990; see also  
174 Ferguson, 1998) are both designed to calculate phase (degrees) and amplitude (dB > 1



175  $\mu\text{V/m}$ ) as functions of distance when supplied with transmitter frequency, radiated  
176 power and ionospheric  $D$  region characteristics,  $H'$  and  $\beta$ , or (particularly for  
177 ModeFinder) an electron density versus height profile. The observed phase and  
178 amplitude, on each path at each time, are then compared with calculated values over  
179 appropriate ranges of  $H'$  and  $\beta$ , to determine the actual  $H'$  and  $\beta$  at that time on that  
180 path, as detailed in sections 3, 4, and 5 below. The  $H'$  and  $\beta$  so determined for each  
181 path at each time are thus their averages along the path for each (constant) SZA.  
182 There will, in reality, be some variation along these paths due to the latitudinal  
183 dependence of galactic cosmic rays as mentioned in the Introduction above and in  
184 Section 6.1 below.

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186 Quite often there is not much to choose between ModeFinder and LWPC; they both  
187 give very similar results under most conditions (e.g., Thomson et al., 2017). Some  
188 considerations in choosing, taken into account here, are briefly discussed below in  
189 section 6.3.

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### 191 **3. Results: Mid-latitude Observations**

#### 192 **3.1 NSY, on Sicily in the Mediterranean, to Cambridge, UK**

193 Figure 2 shows the morning (0-12 UT) phase and amplitude observations (top panels)  
194 recorded at Oakington (52.26°N, 0.07°E), near ( $\sim$ 10 km N.W. of) Cambridge, UK, in  
195 the summer of 2005 (30 July to 4 August) transmitted from the 45.9 kHz transmitter,  
196 NSY (37.13°N, 14.44°E), 2023 km to the southeast as shown in Figure 1a. VLF phase  
197 plots, such as those here, typically need to be corrected (adjusted), for convenience,  
198 for small, fairly constant phase drifts and occasional (often random) phase jumps,  
199 normally at the transmitter, to allow the resulting daily phase plots, over several days,

200 to superpose (approximately) as in Figure 2. Here, for NSY-Oakington, no phase  
201 drift correction was actually needed, and the only phase jumps needing correction  
202 were a very small number involving only multiples of  $180^\circ$  which is common for  
203 MSK modulation (e.g., Thomson, 2010) as used by all the VLF transmitters here. The  
204 middle panels of Figure 2 show the same phase and amplitude observations but with  
205 an expanded time scale about dawn ( $\sim 3.2$ - $4.6$  UT) with  $SZA = 90^\circ$ ,  $97^\circ$ , and  $98^\circ$  being  
206 indicated with vertical dashed lines. The path was fairly well aligned with the sunrise  
207 terminator (see Figure S1a in the Supporting Information, together with Figures S1b,  
208 S1c, S1d ... S1k for all 11 paths); at dawn the  $SZA$  varied by  $\sim 2.5^\circ$  ( $-90^\circ \pm 1.2^\circ$ ) along  
209 the 2023-km path, while near  $SZA \approx 97.5^\circ$ , it varied by only  $\sim 0.4^\circ$  along the path.  
210 Also shown in Figure 2 (bottom panels) are ModeFinder calculations of the phases  
211 and amplitudes of NSY to be expected near Cambridge for a range of appropriate  
212 parameters  $H'$  and  $\beta$  for the  $D$  region. The observational plots in the upper panels  
213 have been calibrated (by vertically shifting each whole plot) so that their mid-day  
214 values match the calculated values in the lower panels for  $H' = 70.8$  km and  $\beta = 0.42$   
215  $\text{km}^{-1}$ ; these  $H'$  and  $\beta$  values were estimated for the NSY-Cambridge path here from  
216 the calibrated VLF observations of Thomson et al. (2011a, 2012). This removes the  
217 need to know the actual power and phase radiated at the NSY transmitter.

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219 The observed phase in degrees together with the observed amplitude in dB at  
220 particular UT times in the upper panels were then used in the bottom panels to find  
221 the values of  $H'$  and  $\beta$  at each of these selected UT times. The  $SZA$  at the mid-point  
222 of the NSY-Oakington path was then determined at that time from either NOAA's  
223 solar zenith angle calculator (<https://gml.noaa.gov/grad/solcalc>) or, in particular from  
224 their spreadsheets, (<https://gml.noaa.gov/grad/solcalc/calcdetails.html>) or from

225 ‘Google Earth’ by measuring the distance in degrees from the path mid-point to the  
226 sub-solar point on the Earth’s surface for the relevant time and date. The resulting  
227 values of  $H'$  and  $\beta$  as functions of SZA are shown plotted in the top, left-hand panels  
228 of Figure 3 (for  $H'$ ) and Figure 4 (for  $\beta$ ) with black lines and black square plot  
229 symbols. The SZA values in degrees are shown as negative before mid-day and  
230 positive after mid-day in this report. Note that at night, on the far left of these plots,  
231  $H' = \sim 85$  km and  $\beta = \sim 0.7$  km<sup>-1</sup>, very similar to previous observations (Thomson et  
232 al., 2007; Thomson & McRae, 2009), while at mid-day,  $H' = \sim 70.8$  km and  $\beta = \sim 0.42$   
233 km<sup>-1</sup> as discussed above. Note, in particular, the clear minimum at  $H' = \sim 72.5$  km near  
234 dawn (SZA  $\approx -90^\circ$ ).

235 As can be seen in the middle panels of Figure 2, both the amplitude and phase of NSY  
236 recorded at Cambridge start responding to the approach of dawn when the SZA at the  
237 mid-point is between  $-98.0^\circ$  and  $-97.5^\circ$ . As mentioned above the variation of SZA  
238 along the path is only  $\sim 0.4^\circ$  in these conditions. In comparison, Reid (1987) reported  
239 the average pre-dawn onset of mesospheric echoes from the 50-MHz radar at Poker  
240 Flat in Alaska occurred around a mean SZA of  $94^\circ$  with only 7% of echoes occurring  
241 with SZA  $> 96^\circ$ . Reid (1987) also cites Sechrist (1968) as reporting VLF observations  
242 (amplitude only) at Wallops Island, Virginia, showing pre-dawn changes at  $94^\circ$ , and  
243 on one occasion at  $98^\circ$ . However, as discussed further below in section 6.5, Thomas &  
244 Harrison (1970) found, using VLF and LF measurements on paths in the UK, that the  
245 pre-dawn changes in the  $D$  region started at SZA  $\sim 98^\circ$ , very similar to the NSY-  
246 Cambridge result found here.

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### 250 3.2 ICV, on Sardinia in the Mediterranean, to Cambridge, UK

251 Figure S2, in the Supporting Information, shows the phase and amplitude  
252 observations (top and middle panels) recorded at Oakington near Cambridge, UK, in  
253 the summer of 2005 (26 June-1 July) transmitted from the 20.27-kHz transmitter, ICV  
254 (40.92°N, 9.73°E), 1459 km approximately to the southeast as shown in Figure 1a.  
255 Also shown in Figure S2 (bottom panels) are ModeFinder calculations of the phases  
256 and amplitudes of ICV to be expected at Oakington for a range of appropriate  $H'$  and  
257  $\beta$  for these conditions. The observations in the upper panels have been calibrated by  
258 vertically shifting so that their mid-day values match the calculated values in the  
259 bottom panels for  $H' = 71.0$  km and  $\beta = 0.40$  km<sup>-1</sup>. These  $H'$  and  $\beta$  values were  
260 estimated for the ICV-Cambridge path here from the calibrated VLF observations of  
261 Thomson et al. (2011a, 2012). As for most of the other transmitter-receiver-pair  
262 recordings used here (including NSY, Sicily, in section 3.1 above), this removes the  
263 need to know the actual power and phase radiated at the transmitter.  
264 As with NSY-Oakington above,  $H'$  and  $\beta$  as functions of SZA were determined and  
265 are plotted in the same (top, left-hand) panels of Figures 3 and 4 used in section 3.1  
266 but with brown lines and brown diamond plot symbols. Clearly both the NSY and  
267 ICV plots are rather similar though the ICV minimum in  $H'$  near dawn is slightly less  
268 deep (~73.5 km as opposed to ~72.5 km). Although both paths pass over the southern  
269 European (Swiss-French-Italian) Alps, the NSY path is longer but has an alpine  
270 section which is much shorter than the ICV path's alpine section (which runs along  
271 the mountainous French-Italian) border. The NSY path results will thus be assumed to  
272 be likely to be more reliable than the ICV results.

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**276 3.3 Central France northwards to Cambridge, UK (~600 km)**

277 The VLF transmitter at 46.71°N, 1.25°E (in central France near Rosnay), is referred

278 here to as 'FRA', as it was in Thomson (1993). Two recorded periods are used here:

279 for the first, 14-20 June 2005, FRA was radiating on 21.75 kHz while for the second,

280 20-26 June 2005, FRA radiated on 18.3 kHz. The path is shown in Figure 1a.

281 Unfortunately FRA was not using a conventional stable phase, but rather its phase

282 was continually (quasi-) randomly scattered. Our receivers could track this and so we

283 could record valid amplitudes but the phase itself was too unpredictable for us to use.

284 This phase issue might have been resolvable but, at the time, we were unable to

285 record the phase close to the transmitter as well as at the receiver at Oakington. Hence

286 we use only amplitude for FRA-Oakington here; thus the observational plots and

287 ModeFinder calculation plots in Figure S3 are for amplitude only. To enable  $\beta$  to be

288 determined,  $H'$  was assumed to vary with SZA as for NSY-Oakington except that in

289 some instances  $H'$  needed to be adjusted slightly, e.g., at mid-day where  $H' = 72.0$  km

290 and  $\beta = 0.38 \text{ km}^{-1}$  was estimated from Thomson et al. (2011a, 2012). The resulting  $H'$

291 and  $\beta$  as functions of SZA are plotted in the same (top, left-hand) panels of Figures 3

292 and 4 used in section 3.1 but using cyan (18.3 kHz) and blue (21.75 kHz) for the lines

293 and plot symbols. While the  $H'$  plots for the two FRA frequencies are necessarily

294 rather similar to the corresponding NSY plot (because both are derived from the NSY

295  $H'$  data), the  $\beta$  plots are rather similar too indicating the FRA data is providing some

296 support through its apparent consistency with the NSY and ICV data.

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### 300 **3.4 NAU, Puerto Rico, Caribbean, to St. John's, NFL, Canada (~3500 km)**

301 Figure S4 shows phase and amplitude observations (upper panels) recorded at St  
302 John's (47.60°N, 52.68°W), Newfoundland, Canada, 11-18 June 2013, from the  
303 40.75kHz transmitter, NAU (18.40°N, 67.18°W), 3497 km to the southwest as shown  
304 in Figure 1b. The data are analysed during the period around dusk in June, when the  
305 SZA varied by only  $\sim 4^\circ$  ( $\sim 90^\circ \pm 2.5^\circ$ ) along the 3497-km path. Also shown in Figure  
306 S4 (bottom panels) are the appropriate ModeFinder calculations of the phases and  
307 amplitudes. Again, the observational plots in the upper panels have been calibrated  
308 (by vertically shifting each whole plot) so that their mid-day values match the  
309 calculated values in the bottom panels for  $H' = 70.5$  km and  $\beta = 0.43$  km<sup>-1</sup>, as in  
310 Thomson et al. (2014).

311 The resulting values of  $H'$  and  $\beta$  as functions of SZA, which have previously been  
312 reported in Thomson et al. (2014), are shown plotted in a slightly different form and  
313 to higher values of SZA, in the top, right-hand panels of Figure 3 (for  $H'$ ) and Figure  
314 4 (for  $\beta$ ) with the black line and black square plot symbols. Note that at night, on the  
315 far right of these plots,  $H' = \sim 84$  km and  $\beta = \sim 0.6$  km<sup>-1</sup>, in line with previous  
316 observations on other paths (Thomson et al., 2007; Thomson & McRae, 2009), while  
317 at mid-day,  $H' = \sim 70.5$  km and  $\beta = \sim 0.43$  km<sup>-1</sup> as discussed above. Note, in particular,  
318 there is no deep minimum near dusk (SZA = 90°), like that of  $H' = \sim 73$  km near dawn  
319 in the top-left panel of Figure 3. The dusk variation of  $H'$  versus SZA shown here is,  
320 in fact, quite similar to that reported by Thomson (1993) from other earlier paths.

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### 324 **3.5 Central France northwards to Cambridge, UK (~600 km)**

325 As in section 3.3 above, we again use only amplitude for FRA-Oakington, using the  
326 plots in the top panels of Figure S3 but for dusk rather than dawn. To enable  $\beta$  to be  
327 determined,  $H'$  at dusk was assumed to vary with SZA as for NAU to St. John's (as in  
328 section 3.4 above) except that again, in some instances,  $H'$  needed to be adjusted  
329 slightly, e.g., at mid-day where again we took  $H' = 72.0$  km and  $\beta = 0.38$  km<sup>-1</sup> for  
330 FRA-Oakington. The resulting  $H'$  and  $\beta$  as functions of SZA are plotted in the same  
331 (top, right-hand) panels of Figures 3 and 4 used in section 3.4 but using cyan (18.3  
332 kHz) and blue (21.75 kHz) for the lines and plot symbols. While the  $H'$  plots for the  
333 two FRA frequencies are necessarily rather similar to the corresponding NAU plot  
334 (because both are derived from the NAU  $H'$ -data), the  $\beta$  plots can be seen to be rather  
335 similar too, indicating the FRA data is again providing some support by being  
336 consistent with the NAU data.

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## 338 **4. Results: Low Latitude Observations**

### 339 **4.1 NPM, Oahu, Hawaii, to Keauhou, Big Island, Hawaii (306 km)**

340 Thomson et al. (2014, figure 4) measured the phase and amplitude of the 21.4 kHz  
341 U.S. Navy transmitter, NPM (21.42°N, 158.15°W) on Oahu, as received, 306 km  
342 away over a nearly all-sea path, at Keauhou (19.58°N, 155.97°W) on the west coast of  
343 the 'Big Island' of Hawaii (Figure 1d), in the period 19-25 August 2012. These  
344 measurements were then used, together with ModeFinder calculations, to determine  
345  $H'$  and  $\beta$  as functions of SZA at this low latitude, ~20°N (geographic and  
346 geomagnetic: [https://www.ncei.noaa.gov/sites/default/files/2022-03/Geomagnetic](https://www.ncei.noaa.gov/sites/default/files/2022-03/Geomagnetic_Coordinates.pdf)  
347 [Coordinates.pdf](https://www.ncei.noaa.gov/sites/default/files/2022-03/Geomagnetic_Coordinates.pdf)). These  $H'$  and  $\beta$  values are plotted as red lines with red diamond plot  
348 symbols in the middle panels of Figures 3 and 4 respectively, on the left for the  
349 morning values and on the right for the afternoon values. Two extra dB/deg data

350 points, one from 16 UT (87 dB, 40°) and the other from 5 UT  $\equiv$  29 UT (86 dB, 54°)  
351 from Thomson et al. (2014, figure 4) were used here with the extended, NPM-  
352 Keauhou, ModeFinder plots in Figure S5 to generate the two extra SZA data points  
353 ( $H' = 83.0$  km,  $\beta = 0.255$  km<sup>-1</sup> and  $H' = 82.5$  km,  $\beta = 0.260$  km<sup>-1</sup>) not shown in  
354 Thomson et al. (2014) but shown in Figures 3 and 4 here. It is immediately clear in  
355 Figure 3 that, for  $H'$ , the clear minimum ( $\sim 73$  km) at dawn at mid-latitudes does not  
356 occur for the low latitude of Hawaii while, in the afternoon and at dusk (top and  
357 middle, right-hand panels), the variation of  $H'$  with SZA shows similarities at both  
358 mid- and low latitudes.

359

#### 360 **4.2 NWC, North West Cape, Australia, to Karratha, Australia (300 km)**

361 As well as reporting VLF measurements of NPM in Hawaii together with the  
362 resulting  $H'$  and  $\beta$  as functions of SZA used in section 4.1 above, Thomson et al.  
363 (2014) also compared these with  $H'$  and  $\beta$  as functions of SZA for the  $\sim 300$ -km, low  
364 latitude path from NWC (21.82°S, 114.17°E) to Karratha (Millars Well, 20.74°S,  
365 116.82°E) in northwest Australia (Figure 1e) in October 2011. Figure S6 (top panels)  
366 shows the underlying phase and amplitude measurements together with (lower panels)  
367 the appropriate ModeFinder calculations used there and further used again here to  
368 extend the  $H'$  and  $\beta$  results to slightly higher SZA's (similar to the small extensions  
369 for NPM in section 4.1 above). More details about these NWC-Karratha observational  
370 techniques are given by Thomson et al. (2012); the techniques were very similar to  
371 those for NPM-Keauhou, above.

372 These  $H'$  and  $\beta$  values are plotted as functions of SZA using brown lines with open  
373 '+' plot symbols in the middle panels of Figures 3 and 4 respectively, on the left for  
374 the morning values and on the right for the afternoon values. It can thus be seen that



375 the  $H'$  plots for NWC-Karratha are rather similar to those for NPM-Keauhou; in  
376 particular neither of these two low latitude plots shows the marked minimum at dawn  
377 which was apparent for the mid-latitude paths (NSY and ICV to Oakington). The  
378 geographic latitude of the NWC-Karratha path ( $\sim 20^\circ$  S) is very similar in magnitude  
379 to the geographic latitude of the NPM-Keauhou path ( $\sim 20^\circ$  N) but the magnitude of  
380 the geomagnetic latitude of the NWC path is quite a bit higher,  $\sim 31^\circ$  S (versus  $20^\circ$  for  
381 the NPM path). Thus the increased GCR at this somewhat higher latitude is not  
382 sufficient to produce a discernable dawn minimum.

383

#### 384 **4.3 NWC, North West Cape, Australia, to Uji, Kyoto, Japan (6,680 km)**

385 Araki et al. (1969) reported the variations with time (JST = UT+9) of the observed  
386 average phase and amplitude of NWC, when radiating on 15.5 kHz, across the  
387 equator (Araki, 1973), and observed at Uji ( $34^\circ 54'N$ ,  $135^\circ 48'E$ ), Kyoto, Japan (Figure  
388 1c), during the period 31 July – 7 August 1968. The path was fairly closely aligned  
389 with the dusk terminator during this period; the alignment with the dawn terminator  
390 was only rather approximate but the absence of clear mode conversion effects near  
391 dawn probably indicates that the alignment will none-the-less suffice.

392 The bottom two panels of Figure S7 show the results of LWPC calculations of phase,  
393 in degrees, and amplitude, in dB  $> 1 \mu V/m$ , for appropriate values of  $H'$  and  $\beta$  for a  
394 radiated power of 1.0 MW at 15.5 kHz. From Thomson et al. (2012), the (average)  
395 mid-day parameters for the  $D$  region on this trans-equatorial path in 1968 (high solar  
396 maximum) were estimated to be  $H' = 69.5$  km and  $\beta = 0.50 \text{ km}^{-1}$ , which, as can be  
397 seen in the LWPC-calculated plots in the (bottom panels) of Figure S7, gives  $29^\circ$  and  
398 66.5 dB at mid-day. These two values were then used to normalize the phase and  
399 amplitude observations in Araki et al. (1969) which are shown, as functions of JST, in

400 the top two panels of Figure S7, thus enabling  $H'$  and  $\beta$  to be determined, from the  
401 bottom panels, for each time (and hence subsequently each SZA), in a similar way to  
402 that in section 3.1. The resulting  $H'$  and  $\beta$  as functions of SZA (many of which were  
403 reported in Thomson et al., 2014) are shown, using violet lines with violet circles as  
404 plot symbols in the middle two panels of Figures 3 and 4 respectively, on the left for  
405 the morning values and on the right for the afternoon values. Again there is no  
406 minimum in  $H'$  near dawn for this fairly low latitude trans-equatorial path as was also  
407 seen for the other low latitude paths above (sections 4.1 and 4.2). The geomagnetic  
408 latitude of Kyoto is also low at about  $26^\circ\text{N}$ .

409

#### 410 **4.4 NPM, Oahu, Hawaii, to Dunedin, NZ (8090 km)**

411 The path from NPM in Hawaii to Dunedin, NZ (Thomson et al., 2011b & 2012) is  
412 mainly low latitude (Figure 1f), passing ~southwest across the equator, but becomes  
413 mid-latitude for the final part from  $\sim 30^\circ\text{S}$  to the Dunedin receiver ( $45.79^\circ\text{S}$ ,  
414  $170.48^\circ\text{E}$ ), for which the geomagnetic latitude is  $\sim 49^\circ\text{S}$

415 ([https://www.ncei.noaa.gov/sites/default/files/2022-03/Geomagnetic](https://www.ncei.noaa.gov/sites/default/files/2022-03/GeomagneticCoordinates.pdf)  
416 [Coordinates.pdf](https://www.ncei.noaa.gov/sites/default/files/2022-03/GeomagneticCoordinates.pdf)).

417 Figure S8 (upper four panels) shows the measured phases and amplitudes (normalized  
418 at mid-day to  $378^\circ$  and 53.2 dB (using  $H' = 70.6$  km and  $\beta = 0.45$   $\text{km}^{-1}$ ), as functions  
419 of time-of-day, in UT, for the period 16-22 January 2017 UT during which the dawn  
420 terminator is fairly closely aligned with this ( $\sim 8.1$  Mm) path. Also shown (two bottom  
421 panels) are LWPC calculations for phase and amplitude at Dunedin for a range of  
422 appropriate values (as averages along the path) of  $H'$  and  $\beta$ , thus enabling specific  
423 values of  $H'$  and  $\beta$ , to be determined for a range of morning values of UT and hence  
424 SZA. These are shown plotted in the middle, left-hand panels of Figures 3 and 4,

425 respectively, from which it can be seen that there is a minimum at dawn ( $SZA=-90^\circ$ ,  
 426  $H' = \sim 76.5$  km), but this minimum is markedly less deep than that ( $H' = \sim 73$  km) for a  
 427 fully mid-latitude path (e.g., NSY-Oakington, top left panel Figure 3), consistent with  
 428 the NPM-Dunedin path being partly at mid-latitudes, though mainly at low latitudes  
 429 where GCR electrons are contributing only modestly to the lowest ionosphere.

430

## 431 **5. Results: High Mid-Latitude Observations ( $\sim 53^\circ$ - $65^\circ$ Geomagnetic)**

### 432 **5.1 DHO, North Germany to Eskdalemuir, Scotland (748 km)**

433 As shown in Figure 1a, VLF transmissions from DHO ( $53.08^\circ\text{N}$ ,  $7.61^\circ\text{E}$ ) after  
 434 travelling 748 km, mainly across the North Sea, were recorded at Eskdalemuir  
 435 ( $55.31^\circ\text{N}$ ,  $3.21^\circ\text{W}$ ) in Scotland, during the period 4-19 July 2015 UT. These  
 436 recordings were used in Thomson et al. (2017) to determine  $H'$  and  $\beta$ , for this high  
 437 mid-latitude path, as functions of  $SZA$  up to  $\sim 75^\circ$ . Here these results are extended to  
 438  $SZA$ 's beyond  $90^\circ$ . Figure S9 (upper four panels) shows the measured phases and  
 439 amplitudes as functions of UT, normalized at mid-day as in Thomson et al. (2017), at  
 440  $120^\circ$  and 68.8 dB using  $H' = 72.8$  km and  $\beta = 0.345$   $\text{km}^{-1}$  in ModeFinder; also shown  
 441 (bottom two panels) are plots of ModeFinder-calculated phases and amplitudes for  
 442 appropriate ranges of  $H'$  and  $\beta$ .

443 The bottom left-hand panels of Figures 3 and 4 show, as green lines and diamond plot  
 444 symbols, the resulting morning values of  $H'$  and  $\beta$  as functions of  $SZA$  for the DHO-  
 445 Eskdalemuir path. As can be seen, this high mid-latitude path also shows a deep  
 446 minimum in  $H'$  at dawn ( $SZA=-90^\circ$ ,  $H' = \sim 72.8$  km), consistent with enhanced  
 447 electron density in the lowest  $D$  region, as the visible light of dawn releases electrons,  
 448 generated overnight by GCR, from their attached negative ions. The bottom right-  
 449 hand panels of Figures 3 and 4 show the resulting afternoon values of  $H'$  and  $\beta$  as

450 functions of SZA for the DHO-Eskdalemuir path. The dawn terminator is fairly well  
451 aligned with the path during the July measurements here; the SZA varies by only  
452  $\sim 3.5^\circ$  along the path at dawn (SZA  $\sim -90^\circ$ ). At dusk it is not well aligned with the dusk  
453 terminator, but the path is short resulting in the SZA at dusk varying only  $\sim 6.5^\circ$  along  
454 its 748-km length. It thus seems likely that mode conversion should not be too  
455 significant, at least below SZA  $\sim +90^\circ$ , before  $\beta$  starts to change more rapidly with  
456 SZA.

457

## 458 **5.2 NAA, Maine, USA, to St. John's, Newfoundland, Canada (1173 km)**

459 As indicated in Figure 1b, VLF transmissions from NAA (44.64°N, 67.28°W), after  
460 travelling over a mixed land-sea path, were recorded at St. John's, Newfoundland  
461 (47.60°N, 52.68°W) during the period 1-28 June 2019. The SZA changes along the  
462 length of this 1173-km path by  $\sim 10^\circ$  near dawn and by  $\sim 7^\circ$  near dusk. The path is  
463  $\sim 50\%$  over the sea and  $\sim 50\%$  over land; this land has low electrical conductivity  
464 (Morgan, 1968; ITU-R, 2015) mainly  $\sim 0.001$  S/m but some parts have  $\sim 0.0003$  S/m  
465 near St. John's (compared with western Europe at mainly  $\sim 0.01$  S/m). Figure S10  
466 (upper four panels) shows the measured phases and amplitudes as functions of UT,  
467 normalized at mid-day as in Thomson et al. (2017), at  $313^\circ$  and 74.8 dB using  $H' =$   
468  $72.0$  km and  $\beta = 0.33$  km $^{-1}$  in LWPC; also shown (lower two panels) are plots of  
469 LWPC-calculated phases and amplitudes for appropriate ranges of  $H'$  and  $\beta$ . The  
470 resulting morning values of  $H'$  and  $\beta$  as functions of SZA for the NAA-St. John's path  
471 are shown in blue (triangular plot symbols) in the left-hand bottom panels of Figures  
472 3 and 4 respectively while the corresponding plots for the afternoon values are shown  
473 in the right-hand bottom panels. Interestingly there are small, marginally significant,  
474 minima at dusk (SZA  $= +90^\circ$ ) in both of the  $H'$  curves at high-mid-latitude (NAA-St.

475 John's and DHO-Eskdalemuir); these small minima may not be significant or may be  
476 due to mode-conversion discussed briefly in section 6.1 below.

477

### 478 **5.3 22.1 kHz from Skelton, North England, to Reykjavik, Iceland (1487 km)**

479 The 22.1-kHz transmitter at 54.73°N, 2.88°W near Skelton, Cumbria, England, is  
480 referred to here as 'GQD' as in Thomson et al. (2018) and in Koh et al. (2019). This  
481 same 22.1-kHz transmitter near Skelton has also been referred to as 'GVT' (e.g.,  
482 Clilverd et al., 2020). The 1478-km path from GQD, Skelton, to Reykjavik (64.11°N,  
483 21.79°W) travels ~northwest over Scotland (Figure 1a) but is ~75% over the sea, with  
484 <~10% of the path over low conductivity ground (~0.001 S/m), and with the rest of  
485 the land with good conductivity (0.01 S/m). Results from this path have been reported  
486 by Thomson et al. (2018) but a wider range of times-of-day and SZA are used here.  
487 Figure S11 shows (upper four panels) the measured phases and amplitudes as  
488 functions of UT, for the periods 19-25 July 2016 and 19-25 August 2016 normalized  
489 at mid-day for the July period at 11.5° and 64.6 dB using  $H' = 73.5\text{km}$  and  $\beta = 0.33$   
490  $\text{km}^{-1}$  in LWPC and for the August period at 14.5° and 64.5 dB using  $H' = 73.5\text{km}$  and  
491  $\beta = 0.32 \text{ km}^{-1}$  in LWPC; also shown (lowest two panels) are plots of LWPC-  
492 calculated phases and amplitudes for appropriate ranges of  $H'$  and  $\beta$ . The resulting  
493 morning values of  $H'$  and  $\beta$ , as functions of SZA for this GQD-Reykjavik path, are  
494 shown in the bottom left-hand panels of Figures 3 and 4 respectively. The dawn  
495 terminator is well aligned with the path during the July period but somewhat less so  
496 for August; the SZA varies by only ~1° along the path at dawn (SZA~90°) in July but  
497 for August varies ~5°.

498

## 499 **6. Discussion**

500 Dawn/dusk propagation conditions have been considered for three separate latitudinal  
501 ranges. The assumption made is that latitudinal differences in the D-region response  
502 to solar EUV and GCR forcing are large in comparison to any longitudinal variations.  
503 However, we note here that the low-latitude analysis presented here uses paths that  
504 are primarily focused on the Pacific region, while the mid and mid-high latitude  
505 analysis relies on paths focused on the North Atlantic region. This experimental  
506 restriction on any longitudinal interpretation is an area for future work, requiring new  
507 datasets to be developed. At mid-high latitudes the balance between geographic  
508 coordinate systems and geomagnetic coordinate systems could be a factor in  
509 considering longitudinal variations in D-region responses to dawn/dusk conditions,  
510 especially in response to the geomagnetic field effect on GCR intensity. An example  
511 of this is shown in Figure 3 of Clilverd et al. (1991) where geomagnetic longitude  
512 contours (in the form of L-shell contours) are over-plotted on a geographic map and  
513 show substantial variations in the Atlantic region compared with the Pacific. Never-  
514 the-less, the analysis presented here clearly shows the importance of SZA on D-region  
515 electron concentration profiles, as is discussed in more detail below.

### 516 **6.1 Dependency of $H'$ on SZA**

517 In Figure 3, aside from the pronounced minima at dawn ( $SZA=-90^\circ$ ) at mid-latitudes,  
518 the variations of  $H'$  with SZA are observed to be broadly quasi-Chapman-like as  
519 reported by Thomson (1993). Away from dawn and dusk,  $D$  region electron density  
520 profiles, at least at low and middle latitudes, are largely generated by 121.6 nm solar  
521 Lyman- $\alpha$  ionizing the minor neutral constituent  $NO \rightarrow NO^+ + e^-$ , but with the  
522 Lyman- $\alpha$  depth penetration being controlled through absorption by the major neutral  
523 constituent  $O_2$  (Banks & Kockarts, 1973). This results in Lyman- $\alpha$  ionization

524 penetrating down to as low as ~70 km altitude at mid-day but to increasingly greater  
525 heights as the SZA increases towards dawn and dusk.

526 Below 65-70 km, ionization is dominated by galactic cosmic rays continuously  
527 (partially) ionizing all neutral constituents. This is essentially offset by the free  
528 electrons so generated rapidly becoming attached to some of the abundant neutral O<sub>2</sub>  
529 molecules forming negative ions,  $O_2 + e^- + M \rightarrow O_2^- + M$ , where M is another  
530 neutral molecule, typically O<sub>2</sub> or N<sub>2</sub>, enabling conservation of both momentum and  
531 energy (Banks & Kockarts, 1973). During the night these, in turn, exchange with  
532 other neutral species forming other (hydrated) negative ions such as CO<sub>3</sub><sup>-</sup> and NO<sub>3</sub><sup>-</sup>  
533 (e.g., Reid, 1976, 1987; Verronen et al., 2006) which, all being so much more massive  
534 than free electrons, have essentially no effect on VLF radio waves. Thus essentially  
535 no free electrons are observed at these low altitudes at night.

536 Photo-detachment at dawn from the CO<sub>3</sub>-based ions requires only visible light while  
537 UV is needed for the NO<sub>3</sub>-based ions. As dawn approaches, the first changes in the  
538 VLF phases and amplitudes are seen when the SZA is between -98° and -97.5°  
539 (section 3.1 above) which is presumably due to visible light, because of the nearly  
540 horizontal incoming sunlight being depleted of UV by low altitude ozone (e.g.,  
541 Macotela et al., 2019). Photo-detachment at dawn from the NO<sub>3</sub>-based ions will  
542 presumably come later. Hence when sunlight returns around dawn, the electrons are  
543 photo-detached from these negative ions accumulated during the night, resulting in  
544 more cosmic ray generated free electrons at these low altitudes than at any other time  
545 of day and is likely to account for the marked minimum in  $H'$  of ~73 km seen at dawn  
546 at mid-latitudes in Figure 3. This is also believed to be the cause of the VLF phase  
547 overshoots often observed at dawn (e.g., Macotela et al., 2019, and references therein;  
548 see also, e.g., top left panels of Figures 2, S2, S8, and S9). Detailed D-region

549 chemistry simulations are required to definitively show that this is the cause. In  
550 contrast, no minimum in  $H'$  at dawn can be seen at low latitudes (middle, left-hand  
551 panel of Figure 3) for the observations of NPM in Hawaii or for the trans-equatorial  
552 path NWC to Kyoto ( $\sim 31^\circ\text{S}$  through  $0^\circ$  to  $26^\circ\text{N}$  geomagnetic), and only a tentative  
553 small minimum (at  $H' \sim 81$  km) can be seen for the short NWC-Karratha path ( $\sim 30^\circ\text{S}$ ).  
554 This is likely to be due to the markedly smaller galactic cosmic ray intensity at these  
555 low latitudes (Neher & Anderson, 1962; Heaps, 1978). This middle, left-hand panel  
556 of Figure 3 also shows a moderate dawn minimum ( $\sim 76.5$  km) for the trans-equatorial  
557 path NPM to Dunedin, NZ ( $\sim 21^\circ\text{N}$  through  $0^\circ$  to  $49^\circ\text{S}$  geomagnetic) for which  $\sim 70\%$   
558 is low latitude and  $\sim 30\%$  is mid-latitude. Further, all the high mid-latitude paths in  
559 Figure 3 (bottom, left-hand panel) show the marked dawn minimum (72-73 km). In  
560 contrast, at dusk, in Figure 3 (right-hand panels) there are no minima in  $H'$  or only  
561 relatively minor (even tentative as in the top panel) minima (near  $\text{SZA}=90^\circ$ ). These  
562 latter minor minima may not be significant or may possibly be due to mode-  
563 conversion issues caused by rapidly changing  $H'$  and  $\beta$  near and after dusk. VLF  
564 mode-conversion (e.g., Pappert & Snyder, 1972) occurs when some property (e.g.,  $H'$ ,  
565  $\beta$ , ground conductivity etc.) of the Earth-ionosphere waveguide changes over a short  
566 distance  $< \sim \lambda$  (a wavelength) causing some power to scatter into one or more new  
567 modes resulting in interference between modes.  
568 Figure 3 (bottom panels) also shows that, away from dawn and dusk, variations of  $H'$   
569 with daytime solar zenith angle tend to be smaller at higher path latitudes probably  
570 mainly due to the greater influence of (SZA-independent) galactic cosmic ray  
571 electrons at these higher geomagnetic latitudes. At night, shortly before dawn, the  $H'$   
572 values reported here are mainly 84-85 km consistent with the nighttime  $H'$   
573 observations of Thomson et al. (2007) and Thomson & McRae (2009). Shortly after



574 dusk they tend to be a little ( $\sim 1$  km) lower probably consistent with the  $D$  region  
575 taking up to  $\sim 1$ -2 hours or so to settle into night conditions.

576

## 577 **6.2 Dependency of $\beta$ on SZA**

578 Apart from near dawn and dusk,  $\beta$  shows the clear trend, in Figure 4, of decreasing  
579 with increasing SZA from a maximum near noon. This is due to the higher altitude  
580 Lyman- $\alpha$  electrons decreasing as the Sun gets lower in the sky while the lower  
581 altitude galactic cosmic ray electrons remain constant, being largely SZA independent  
582 (Thomson & Clilverd, 2001; Thomson et al., 2014). At dawn and dusk, as can be  
583 seen,  $\beta$  changes generally fairly rapidly with SZA, particularly at dawn, from  
584 nighttime values of typically  $\sim 0.6$ - $0.8$  km $^{-1}$  to high-SZA daytime values of  $\sim 0.25$ - $0.3$   
585 km $^{-1}$ . Not all paths show nighttime  $\beta$  values of  $\sim 0.6$ - $0.8$  km $^{-1}$ . There are a variety of  
586 reasons for this (which may also apply for nighttime values of  $H'$  here on these paths).  
587 In the case of the short path NPM-Keauhou, no automated recording was available  
588 and it was not convenient to continue the manual recordings at night. For the short  
589 path NWC-Karratha the situation was somewhat similar compounded, as discussed in  
590 Thomson et al. (2014), by low sky-wave amplitude and low ground conductivity.  
591 Trans-equatorial paths, such as NWC-Kyoto and NPM-Dunedin used here, have been  
592 found to be frequently anomalous at night though not by day nor, as found here,  
593 largely not during dawn/dusk transitions. As noted above, the path ICV-Cambridge  
594 (section 3.2) with significant parts over mountain ranges and the path NAA-St. John's  
595 (section 5.2) largely over low conductivity ground are also liable to have nighttime  
596 anomalies (Thomson et al., 2007; Thomson & McRae, 2009). None-the-less, the paths  
597 NSY and FRA to Cambridge, NAU to St. John's, DHO to Eskdalemuir and GQD to

598 Reykjavik (in August) here provided regular nighttime parameters (even though the  
599 night was fairly short for the last two).

600

### 601 **6.3 Early US Navy Recommendations and Waveguide Codes**

602 The US Navy's LWPC propagation code for VLF contains an optional 'LWPM'  
603 ionospheric model which has some limited SZA-dependence (Ferguson, 1998). In this  
604 model,  $H'$  and  $\beta$  are constants by day at 74 km and  $0.3 \text{ km}^{-1}$ , respectively, for all  
605 daytime SZA, unless SZA becomes  $>90^\circ$  when,  $H'$  and  $\beta$  both increase stepwise with  
606 SZA ( $90^\circ$  through  $99^\circ$ ) with  $H'$  increasing from 74 km to 87 km (at night), and  $\beta$   
607 increasing from  $0.3 \text{ km}^{-1}$  to  $0.5 \text{ km}^{-1}$  (at night).  $H'$  and  $\beta$  then maintain these constant  
608 nighttime values of 87 km and  $0.5 \text{ km}^{-1}$ , respectively, for all nighttime SZA.

609 This internal ('LWPM') ionospheric model was not used here because it is outdated.

610 We used LWPC or ModeFinder only with user supplied ( $D$  region) ionospheric data

611 (either as  $H'$  and  $\beta$  or as electron density versus height profiles). LWPC has its own

612 internal conductivity map and so is to be preferred over ModeFinder when there is

613 significant low conductivity ground, particularly if this varies along the path, such as

614 for NAA to St. John's here. LWPC may be preferable, on occasion, because it

615 calculates the geomagnetic field parameters in each segment along the path.

616 ModeFinder is also more suitable for comparing with loop antenna (i.e., magnetic

617 field) observations on short ( $\sim 300$  km) paths (Thomson, 2010; see also Thomson et al.

618 (2014 & 2017). ModeFinder appears to be more stable at higher VLF/LF frequencies,

619 appearing to be less likely to miss modes; it can also be simpler to use, particularly

620 with electron density versus height profiles. Most of these issues were found to be not

621 very significant for the accuracies required here.

622 LWPC was typically not needed (nor preferred to ModeFinder) for short paths or  
623 those that varied slowly along their lengths such as the terminator-aligned paths here.  
624 However, a mode-conversion code such as LWPC will definitely be required for  
625 calculations on paths across a dawn/dusk terminator, using our new rapidly changing  
626 (with SZA) path parameters shown in Figures 3 and 4, unless the path is very short  
627 (~few hundred km) or the path is aligned within a few degrees of the terminator.

628

#### 629 **6.4 Rocket Comparisons**

630 Mechtly & Smith (1970) reported rocket observations near dawn at Wallops Island,  
631 Virginia, USA (38°N, 70°W). The electron number densities were measured as  
632 functions of height using a combination of probe currents and radio propagation  
633 (Faraday rotation together with differential absorption). Data from their dawn profile  
634 taken at SZA=90° on 24 July 1968 is reproduced here in Figure 5 (green circles).  
635 Using this profile in ModeFinder for the NSY-Oakington path (as described above in  
636 section 3.1) results in a calculated phase at the receiver of  $-106^\circ \equiv 254^\circ$  (modulo 360°)  
637 which is 10 ° higher than the observed phase (at SZA=90°) of 244°. When the rocket  
638 profile is raised by 2 km (shown with blue circles in Figure 5), the ModeFinder  
639 calculated phase at the receiver became  $-120^\circ \equiv 240^\circ$ , slightly lower (by 4 °) than the  
640 observed phase. However the calculated amplitudes, 56.2 dB for the actual rocket  
641 profile and 59.3 dB for the raised profile, were significantly higher than the 54.3 dB  
642 actually observed. These higher amplitude are likely due to the rocket techniques not  
643 adequately measuring the electron densities below 60-65 km (e.g. Thomson et al.,  
644 2022).

645 As shown in our Figures 3 and 4, the VLF technique used here on the NSY-  
646 Cambridge path gives  $H' = 72.75$  km and  $\beta = 0.378$  km<sup>-1</sup> for SZA = -90.2° at dawn

647 (and  $H' = 72.2$  km and  $\beta = 0.364$  km<sup>-1</sup> for SZA = -89.2°). Figure 5 shows, as the solid  
648 black line, the corresponding electron number densities from ModeFinder, using  $H' =$   
649  $72.75$  km and  $\beta = 0.378$  km<sup>-1</sup>, resulting in the observed amplitude of 54.3 dB and  
650 phase of -116° at the Cambridge receiver. Clearly this VLF profile matches well with  
651 the height of the Mechtly & Smith (1970) rocket profiles but has more electrons  
652 below ~60 km; these low altitude electrons are immersed in the higher density neutral  
653 atmosphere there and so have high (electron-neutral) collision frequencies, resulting  
654 in the higher VLF attenuation as actually observed. VLF propagation by day is  
655 generally not sensitive to electron densities above height,  $H'$  (e.g., Thomson et al.,  
656 2022, and references therein); in Figure 5 the orange line at 71 km is indicating the  
657 greatest height at which the VLF propagation on this NSY-Cambridge path was found  
658 here (using ModeFinder) to be sensitive to the electron number density.

659

660 Also shown in Figure 5 for further comparison is the rocket-derived FIRI-2018  
661 (Friedrich et al., 2018) electron density profile for an SZA of 90° for July at 45°  
662 latitude (and solar activity F10.7 = 130 sfu though this is non-critical). As can be  
663 seen, the FIRI electron densities are somewhat lower than the others consistent with  
664 the comparisons of Thomson et al. (2022) who found generally good agreement  
665 between VLF-derived electron density profiles and FIRI-2018 at their common  
666 altitudes except at the lowest altitudes at (high) mid-latitudes, due to the lower  
667 sensitivity of the rocket technique at the lowest altitudes and the significant galactic  
668 cosmic ray generation of electrons there.

669

670 **6.5 Early Ground-Based Radio Observations in the UK and Australia**

671 Thomas & Harrison (1970) used VLF and LF measurements (typically reflection and  
672 polarization coefficients) recorded in the UK in the summer months of 1948-1950,  
673 together with full-wave calculations similar to those of Pitteway (1965) to calculate  
674 pre-sunrise electron density versus height profiles over a range of SZA ( $>105^\circ$ ,  $99^\circ$ ,  
675  $98^\circ$ ,  $97^\circ$ ,  $95^\circ$ ,  $93^\circ$ ,  $91.5^\circ$ , and  $90^\circ$ ). The first two of these profiles (SZA  $>105^\circ$  &  $99^\circ$ ,  
676 i.e., essentially nighttime) agree quite well with the much more recent nighttime  
677 rocket-based FIRI (Friedrich et al., 2018) and VLF-based profiles presented in  
678 Thomson et al. (2022). As the SZA decreases from  $97^\circ$  to  $90^\circ$  the Thomas & Harrison  
679 (1970) profiles descend from  $\sim 80$  km to  $\sim 60$  km with the peak electron densities  
680 increasing from  $\sim 2 \times 10^7 \text{ m}^{-3}$  to  $\sim 1.4 \times 10^8 \text{ m}^{-3}$  (at 65 km at  $90^\circ$ ). To compare with our  
681 more recent results, their lowest profile (i.e., for SZA =  $90^\circ$ ) was entered into  
682 ModeFinder for our, similar-latitude,  $\sim 748$ -km path, DHO (Germany) to Eskdalemuir  
683 (Scotland) resulting in 68.5 dB and  $204^\circ$  for the calculated amplitude and phase at the  
684 receiver. While the amplitude is quite similar to that actually observed, their profile  
685 phase of  $204^\circ$  is quite a lot higher than the  $\sim 120^\circ$  actually observed (Figure S9, in our  
686 Supporting Information) which means their dawn (SZA =  $90^\circ$ ) profile is likely to be  
687 too low in height (by  $\sim 4$  km). Also, by comparing with our Figure 5 (where our  $H'/\beta$   
688 profile for NSY-Oakington there is very similar to that for DHO-Eskdalemuir), it can  
689 be seen their electron density at  $\sim 65$  km appears to be too high by a factor of  $\sim 3$ .  
690 None-the-less their results are impressive for such early observations.

691 Smith et al. (1967) determined *D* region electron densities above Armidale, NSW,  
692 Australia ( $30.5^\circ\text{S}$ ,  $151.5^\circ\text{E}$  geographic,  $\sim 36^\circ\text{S}$  geomagnetic) during  $\sim 1963$ -65 using a  
693 medium frequency ( $\sim 1$ -2 MHz) pulsed cross-modulation technique, from night  
694 through dawn to mid-day. Their electron number densities from night to dawn  
695 (SZA= $-90^\circ$ ) are very broadly similar to, though somewhat less than, the UK values of

696 Thomas & Harrison (1970). This is consistent with the Armidale results being at a  
697 lower geomagnetic latitude than those in the UK ( $\sim 36^\circ$  versus  $\sim 55^\circ$ ), resulting in  
698 lower cosmic ray intensity there. The Armidale peak electron density at  $SZA=-90^\circ$   
699 (dawn) at a height of 65 km was  $\sim 8 \times 10^7 \text{ m}^{-3}$ , closer to our  $\sim 7 \times 10^7 \text{ m}^{-3}$  (Figure 5,  
700 65 km) than the UK value  $\sim 14 \times 10^7 \text{ m}^{-3}$  of Thomas & Harrison (1970). When the  
701 ( $SZA=-90^\circ$ ) Armidale electron densities are used in ModeFinder for our NSY-  
702 Oakington path, the predicted phase and amplitude,  $-107^\circ$  and 55.2 dB, are fairly  
703 similar to the observed values,  $-117^\circ$  and 54.4 dB. Most of the phase difference can be  
704 removed by raising the Armidale profile by just over 1 km, while the higher  
705 amplitude for the Armidale profile is probably due to it having too few electrons  
706 below  $\sim 60$  km. However, the geomagnetic latitude of the mid-point of our NSY-  
707 Oakington path ( $\sim 48^\circ\text{N}$ ) is somewhat higher than that of Armidale ( $\sim 36^\circ\text{N}$ ). None-the-  
708 less the very broad agreement between these earlier (UK and Australian) observations  
709 and those reported here seems quite reasonable.

710

## 711 **7. Summary and Conclusions**

712 Characteristics of VLF radio propagation conditions near dawn and dusk in the lower  
713 *D* region (altitudes  $\sim 55$ -90 km) of the Earth's ionosphere were determined using  
714 phases and amplitudes of subionospheric VLF radio waves. The VLF waves analysed  
715 propagated along either fairly short paths ( $\sim 300$ -1000 km) or along paths which  
716 aligned with the dawn/dusk terminator (to keep the *SZA* fairly constant along the  
717 path). For most paths the phases and amplitudes observed at the receiver near (say)  
718 dawn (or dusk) were measured relative to those observed at mid-day on the same day;  
719 these differences were then added to the corresponding phase or amplitude calculated  
720 for mid-day from previously determined *D* region parameters to get dawn/dusk values

721 independent of the phase or amplitude at the transmitter. Such mid-day calculations  
722 are likely to be more accurate at low latitudes or in summer (or at least, not in winter  
723 when the  $D$  region is more variable).

724 The VLF propagation conditions were characterized with the two ‘Wait’ parameters,  
725  $H'$  as a measure of the height and  $\beta$  as a measure of the sharpness (or slope) of the  
726 electron densities in the lower  $D$  region. The work reported here enabled these two  
727 parameters to be determined as functions of SZA in 3 latitude ranges: mid-latitudes,  
728 low latitudes and high mid-latitudes. For the first time these parameters have been  
729 able to be determined through dawn and dusk which should not only provide  
730 observational data for modeling the  $D$  region at these times but also extend the  
731 coverage range of VLF propagation monitoring for geophysical phenomena such as  
732 lightning, solar flares, and energetic particle precipitation.

733 The plots of  $H'$  and  $\beta$  against SZA also illustrate geophysical effects in the  $D$  region  
734 itself, including the greater role (and so SZA-dependence) of solar EUV (Lyman- $\alpha$ ) at  
735 low and mid-latitudes and the greater role of galactic cosmic rays (and so  
736 SZA-independence) at increasingly higher mid-latitudes. In particular, the marked  
737 minimum observed in  $H'$  versus SZA at dawn at mid-latitudes (due to the release of  
738 cosmic ray generated electrons accumulated over night) reported here for the first  
739 time, shows how very different dawn propagation conditions are from dusk at mid-  
740 latitudes. Again, the observed absence of such a minimum at low latitudes emphasizes  
741 the reduced cosmic ray activity there compared with mid- and higher latitudes and  
742 also emphasizes that the dawn day-night transition in the  $D$  region is very different at  
743 low latitudes and mid-latitudes.

744

## 745 **Data Availability Statement**

746 The data underlying our VLF observations reported here are available at  
747 <http://doi.org/10.5281/zenodo.14210954> or (in some cases) from Thomson et al.  
748 (2007, 2012, 2014, 2017).

749 The FIRI-2018 model profiles are available from  
750 <https://figshare.com/s/357cb03b3e5bed649bbc> (Friedrich et al., 2018) or  
751 <https://figshare.com/search?q=FIRI-2018>

752 US Navy code LWPC is available at <https://github.com/mlhutchins/LWPC>

753 The US Navy computer program referred to here as ModeFinder is a slightly modified  
754 version of MODEFNDR (e.g., Thomson, 1993; Nunn & Strangeways, 2000) and  
755 MODESRCH described and listed in Morfitt & Shellman (1976).

756 Solar zenith angles were determined from NOAA's solar calculator at  
757 <https://gml.noaa.gov/grad/solcalc/>, or, often more conveniently and appropriately for  
758 this study, from their spreadsheets at  
759 (<https://gml.noaa.gov/grad/solcalc/calcdetails.html>).

760 Geomagnetic latitudes came from the NOAA/NCEI World Magnetic Model:

761 <https://www.ncei.noaa.gov/products/world-magnetic-model> specifically from:

762 [https://www.ncei.noaa.gov/sites/default/files/2022-03/Geomagnetic Coordinates.pdf](https://www.ncei.noaa.gov/sites/default/files/2022-03/Geomagnetic%20Coordinates.pdf)

763

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768

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770 **References**

771

772 Araki, T., Kitayama, S., & Kato, S. (1969). Transequatorial reception of VLF radio  
773 waves from Australia. *Radio Science*, 4(4), 367-369.774 <https://doi.org/10.1029/RS004i004p00367>

775

776 Araki, T. (1973). Anomalous diurnal changes of transequatorial VLF radio waves.  
777 *Journal of Atmospheric and Terrestrial Physics*, 35, 693-703.778 [https://doi.org/10.1016/0021-9169\(73\)90200-6](https://doi.org/10.1016/0021-9169(73)90200-6)

779

780 Banks, P. M., & Kockarts, G. (1973). *Aeronomy*, New York, NY: Academic.


781

782 Barman, K., Das, B., Pal, S., Haldar, P. K., Midya, S. K., Pal, S., & Mondal, S. K.  
783 (2024). D-region ionospheric disturbances due to the December 2019 solar eclipse  
784 observed using multi-station VLF radio network. *Advances in Space Research* 74(3),  
785 1460-1470. <https://doi.org/10.1016/j.asr.2024.04.049>

786

787 Basak, T., Hobara, Y., Pal, S., Nakamura, T., Izutsu, J., & Minatohara, T. (2024).  
788 Modeling of Solar Eclipse effects on the sub-ionospheric VLF/LF signals observed  
789 by multiple stations over Japan. *Advances in Space Research* 73, 736-746.790 <https://doi.org/10.1016/j.asr.2023.09.063>

791

792 Chand, A. E., & Kumar, S. (2017). VLF modal interference distance and night- time D  
793 region VLF reflection height for west-east and east-west propagation paths to Fiji,  
794 *Radio Science*, 52(8), 1004–1015. <https://doi.org/10.1002/2016RS006221>795 796 Chand, A. E. & Kumar, S. (2021). Earth-Ionosphere Waveguide Model Parameters  
797 Using VLF Transmissions Received in the South Pacific Region. *IEEE Access* 9,  
798 56653-56663. doi:10.1109/ACCESS.2021.3072133799 <https://ieeexplore.ieee.org/document/9399413>

800

801 Cheng, Z., Cummer, S. A., Baker, D. N., & Kanekal, S. G. (2006). Nighttime D region  
802 electron density profiles and variabilities inferred from broadband measurements  
803 using VLF radio emissions from lightning. *Journal of Geophysical Research*,  
804 111(A5). <https://doi.org/10.1029/2005JA011308>

805

806 Chowdhury, S., Kundu, S., Basak, T., Ghosh, S., Hayakawa, M., Chakraborty, S.,  
807 Chakrabarti, S. K., & Sasmal, S. (2021). Numerical simulation of lower ionospheric  
808 reflection parameters by using International Reference Ionosphere (IRI) model and  
809 validation with Very Low Frequency (VLF) radio signal characteristics. *Advances in*  
810 *Space Research*, 67, 1599-1611. <https://doi.org/10.1016/j.asr.2020.12.017>

811

812 Clilverd, M.A., Smith, A. J., & Thomson, N.R. (1991). The annual variation in quiet  
813 time plasmaspheric electron density, determined from whistler mode group delays.  
814 *Planetary and Space Science*, 39 (7), 1059-1067, [https://doi.org/10.1016/0032-](https://doi.org/10.1016/0032-0633(91)90113-O)  
815 [0633\(91\)90113-O](https://doi.org/10.1016/0032-0633(91)90113-O).

816

817 Clilverd, M. A., Rodger, C. J., Thomson, N. R., Brundell, J. B., Ulich, Th.,  
818 Lichtenberger, J., Cobbett, N., Collier, A. B., Menk, F. W., Seppälä, A., Verronen, P.  
819 T., & Turunen E. (2009). Remote sensing space weather events: Antarctic-Arctic

- 820 Radiation-Belt (Dynamic) Deposition-VLF Atmospheric Research Konsortium  
821 network, *Space Weather*, 7(4), S04001. <https://doi.org/10.1029/2008SW000412>  
822
- 823 Clilverd, M. A., Rodger, C. J., Gamble, R. J., Ulich, T., Raita, T., Seppälä, A., Green, J.  
824 C., Thomson, N. R., Sauvaud, J.-A., & Parrot, M. (2010). Ground-based estimates of  
825 outer radiation belt energetic electron precipitation fluxes into the atmosphere.  
826 *Journal of Geophysical Research: Space Physics*, 115(A12).  
827 <https://doi.org/10.1029/2010JA015638>  
828
- 829 Clilverd, M. A., Rodger, C. J., van de Kamp, M., & Verronen, P. T. (2020). Electron  
830 precipitation from the outer radiation belt during the St. Patrick's day storm 2015:  
831 Observations, modeling, and validation. *Journal of Geophysical Research: Space*  
832 *Physics*, 125(2). <https://doi.org/10.1029/2019JA027725>  
833
- 834 Crombie, D. D. (1964). Periodic fading of VLF signals received over long paths during  
835 sunrise and sunset. *Radio Science Journal of Research NBS/USNC- URSI 68D(1)*,  
836 27-34. <http://dx.doi.org/10.6028/jres.068D.012>  
837 [https://nvlpubs.nist.gov/nistpubs/jres/68D/jresv68Dn1p27\\_A1b.pdf](https://nvlpubs.nist.gov/nistpubs/jres/68D/jresv68Dn1p27_A1b.pdf)  
838
- 839 Cummer, S. A., Inan, U. S., & Bell, T. F. (1998). Ionospheric *D* region remote sensing  
840 using VLF radio atmospherics. *Radio Science*, 33(6), 1781-1792.  
841 <https://doi.org/10.1029/98RS02381>  
842
- 843 Ferguson, J. A., & Snyder, F. P. (1990). *Computer programs for assessment of long*  
844 *wavelength radio communications, version 1.0: Full FORTRAN code user's guide,*  
845 *Naval Ocean Systems Center Tech. Doc. 1773, DTIC AD-B144 839, Def. Tech. Inf.*  
846 *Cent.*, Alexandria, Va.  
847
- 848 Ferguson, J. A. (1998). *Computer programs for assessment of long-wavelength radio*  
849 *communications, version 2.0: User's guide and source files.* SPAWAR Technical  
850 Document 3030, San Diego, CA: Space and Naval Warfare Systems Center.  
851 <https://apps.dtic.mil/sti/pdfs/ADA350375.pdf>  
852
- 853 Fishman G. J. & Inan U. S. (1988). Observation of an ionospheric disturbance caused  
854 by a gamma-ray burst. *Nature* 331, 418-420. <https://doi.org/10.1038/331418a0>  
855
- 856 Friedrich, M., Pock, C., & Torkar, K. (2018). FIRI-2018, an updated empirical model of  
857 the lower ionosphere. *Journal of Geophysical Research: Space Physics*, 123(8),  
858 6737–6751. <https://doi.org/10.1029/2018JA025437>  
859
- 860 Funke, B., Dudok de Wit, T., Ermolli, I., Haberreiter, M., Kinnison, D., Marsh, D.,  
861 Nesse, H., Seppälä, A., Sinnhuber, M., & Usoskin, I. (2024). Towards the definition  
862 of a solar forcing dataset for CMIP7, *Geoscientific Model Development*, 17(3),  
863 1217–1227. <https://doi.org/10.5194/gmd-17-1217-2024>  
864
- 865 Han, F. & Cummer, S. A. (2010). Midlatitude daytime *D* region ionosphere variations  
866 measured from radio atmospherics, *Journal of Geophysical Research*, 115, A10,  
867 <https://doi.org/10.1029/2010JA015715>  
868

- 869 Heaps, M. G. (1978). Parameterization of the cosmic ray ion-pair production rate above  
870 18 km. *Planetary and Space Science*, 26, 513-517. [https://doi.org/10.1016/0032-](https://doi.org/10.1016/0032-0633(78)90041-7)  
871 [0633\(78\)90041-7](https://doi.org/10.1016/0032-0633(78)90041-7)  
872
- 873 ITU-R (2015), International Telecommunications Union, Radio Communications Sector,  
874 Recommendation ITU-R P.832-4, World atlas of ground conductivities, P-series,  
875 Radiowave Propagation. [https://www.itu.int/dms\\_](https://www.itu.int/dms_pubrec/itu-r/rec/p/R-REC-P.832-4-201507-I!!PDF-E.pdf) [pubrec/itu-r/rec/p/R-REC-P.832-4-](https://www.itu.int/dms_pubrec/itu-r/rec/p/R-REC-P.832-4-201507-I!!PDF-E.pdf)  
876 [201507-I!!PDF-E.pdf](https://www.itu.int/dms_pubrec/itu-r/rec/p/R-REC-P.832-4-201507-I!!PDF-E.pdf)  
877
- 878 Koh, K., Bennett, A., Ghilain, S., Liu, Z., Pedeboy, S., Peverell, A., & Füllekrug, M.  
879 (2019). Lower ionospheric conductivity modification above a thunderstorm updraught.  
880 *Journal of Geophysical Research: Space Physics*, 124, 6938–6949.  
881 <https://doi.org/10.1029/2019JA026863>  
882
- 883 Kumar, A., & Kumar, S. (2020). Ionospheric D region parameters obtained using VLF  
884 measurements in the South Pacific region. *Journal of Geophysical Research: Space*  
885 *Physics*, 125, e2019JA027536. [https://doi.org/ 10.1029/2019JA027536](https://doi.org/10.1029/2019JA027536)  
886
- 887 McCormick, J. C., & Cohen, M. B. (2021). A new four-parameter D-region  
888 ionospheric model: Inferences from lightning-emitted VLF signals. *Journal of*  
889 *Geophysical Research: Space Physics*, 126, e2021JA029849.  
890 <https://doi.org/10.1029/2021JA029849>  
891
- 892 Macotela, E. L., Clilverd, M. A., Manninen, J., Thomson, N. R., Newnham, D. A., &  
893 Raita, T. (2019) The effect of ozone shadowing on the D region ionosphere during  
894 sunrise, *Journal of Geophysical Research*, 124(5), 3729-3742.  
895 <https://doi.org/10.1029/2018JA026415>  
896
- 897 Macotela, E. L., Clilverd, M., Renkwitz, T., Chau, J., Manninen, J., & Banyás<sup>[SEP]</sup>D.  
898 (2021). Spring-fall asymmetry<sup>[SEP]</sup> in VLF amplitudes recorded in the North Atlantic  
899 region: The fall-effect. *Geophysical Research Letters*, 48, e2021GL094581.  
900 <https://doi.org/10.1029/2021GL094581>  
901
- 902 Matthes, K., Funke, B., Andersson, M. E., Barnard, L., Beer, J., Charbonneau, P.,  
903 Clilverd, M. A., Dudok de Wit, T., Haberreiter, M., Hendry, A., Jackman, C. H.,  
904 Kretzschmar, M., Kruschke, T., Kunze, M., Langematz, U., Marsh, D. R.,  
905 Maycock, A. C., Misios, S., Rodger, C. J., Scaife, A. A., Seppälä, A., Shangguan,  
906 M., Sinnhuber, M., Tourpali, K., Usoskin, I., van de Kamp, M., Verronen, P. T., &  
907 Versick, S. (2017). Solar forcing for CMIP6 (v3.2), *Geoscientific Model*  
908 *Development*, 10(6), 2247–2302. <https://doi.org/10.5194/gmd-10-2247-2017>  
909
- 910 McRae, W. M., & Thomson, N. R. (2004). Solar flare induced ionospheric D-region  
911 enhancements from VLF phase and amplitude observations, *Journal of Atmospheric*  
912 *and Solar-Terrestrial Physics*, 66(1), 77-87.  
913 <https://doi.org/10.1016/j.jastp.2003.09.009>  
914
- 915 Mechtly, E. A., & Smith, L. G. (1970). Changes of lower ionosphere electron  
916 densities with solar zenith angle. *Radio Science*, 5(12), 1407-1412.  
917 <https://doi.org/10.1029/RS005i012p01407>  
918

- 919 Mechtly, E. A., Bowhill, S. A., & Smith, L. G. (1972). Changes of lower ionosphere  
920 electron concentrations with solar activity. *Journal of Atmospheric and Terrestrial*  
921 *Physics*, 34(11), 1899–1907. [https://doi.org/10.1016/0021-9169\(72\)90065-7](https://doi.org/10.1016/0021-9169(72)90065-7)  
922
- 923 Morfitt, D. G., & Shellman, C. H. (1976). “MODESRCH”, an Improved Computer  
924 Program For Obtaining ELF/VLF/LF Mode Constants in an Earth Ionosphere  
925 Waveguide (Naval Electr. Lab. Cent. Interim Rep. 77T, NTIS Accession  
926 ADA032573). Springfield, VA: National Technical Information Service.  
927 <https://apps.dtic.mil/sti/pdfs/ADA032573.pdf>  
928
- 929 Morgan, R. R. (1968). Westinghouse world-wide VLF conductivity map.  
930 AD0675771. <https://apps.dtic.mil/sti/pdfs/AD0675771.pdf>  
931
- 932 Neal, J. J., Rodger, C. J., Clilverd, M. A., Thomson, N. R., Raita, T. & Ulich, Th.  
933 (2015). Long-term determination of energetic electron precipitation into the  
934 atmosphere from AARDDVARK subionospheric VLF observations. *Journal of*  
935 *Geophysical Research*, 120(3), 2194-2211. <https://doi.org/10.1002/2014JA020689>  
936
- 937 Neher, H. V., & Anderson, H. R. (1962). Cosmic rays at balloon altitudes and the  
938 solar cycle, *Journal of Geophysical Research*, 67(4), 1309-1315.  
939 <https://doi.org/10.1029/JZ067i004p01309>  
940
- 941 Nunn, D., & Strangeways, H. J. (2000). Trimpi perturbations from large ionisation  
942 enhancement patches (LIEs), *Journal of Atmospheric and Terrestrial Physics*,  
943 62(3), 189–206. [https://doi.org/10.1016/S1364-6826\(00\)00004-3](https://doi.org/10.1016/S1364-6826(00)00004-3)  
944
- 945 Pal, S., Hobara, Y., Shvets, A., Schnoor, P. W., Hayakawa, M., and Koloskov, O.  
946 (2023). First detection of global ionospheric disturbances associated with the most  
947 powerful gamma ray burst GRB221009A. *Atmosphere*, 14, 217.  
948 <https://doi.org/10.3390/atmos14020217>  
949
- 950 Pappert, R. A., & Snyder, F. P. (1972). Some results of a mode-conversion program  
951 for VLF. *Radio Science*, 7(10), 913-923.  
952 <https://doi.org/10.1029/RS007i010p00913>  
953
- 954 Peterson, J. R. (1976). Sunlight photodestruction of  $\text{CO}_3^-$ ,  $\text{CO}_3^- \cdot \text{H}_2\text{O}$ , and  $\text{O}_3^-$ : The  
955 importance of photodissociation to the *D* region electron densities at sunrise.  
956 *Journal of Geophysical Research*, 81(7), 1433-1435.  
957 <https://doi.org/10.1029/JA081i007p01433>  
958
- 959 Pitteway, M. L. V. (1965). The numerical calculation of wave-fields, reflexion  
960 coefficients and polarizations for long radio waves in the lower ionosphere. I.,  
961 *Philosophical Transactions of the Royal Society*, A257(1079), 219-241.  
962 <https://doi.org/10.1098/rsta.1965.0004>  
963
- 964 Reid, G. C. (1976). Ion chemistry in the *D* region. *Advances in Atomic and Molecular*  
965 *Physics (Edited by Bates, D. R. & Bederson, B.)*, 12, 375-413. Academic Press,  
966 New York. [https://doi.org/10.1016/S0065-2199\(08\)60047-0](https://doi.org/10.1016/S0065-2199(08)60047-0)  
967

- 968 Reid, G. C. (1987). Radar observations of negative-ion photodetachment at sunrise in  
 969 the auroral-zone mesosphere. *Planetary and Space Science*, 35(1), 27–37.  
 970 [https://doi.org/10.1016/0032-0633\(87\)90141-3](https://doi.org/10.1016/0032-0633(87)90141-3)  
 971
- 972 Rodger, C. J., Brundell, J. B., & Dowden, R. L. (2005). Location accuracy of VLF  
 973 World Wide Lightning Location (WWLL) network: Post-algorithm upgrade.  
 974 *Annales Geophysicae*, 23, 277-290. <https://doi.org/10.5194/angeo-23-277-2005>  
 975
- 976 Rodger, C. J., Clilverd, M. A., Seppälä, A., Thomson, N. R., Gamble, R. J., Parrot,  
 977 M., Sauvaud, J.-A., et al. (2010). Radiation belt electron precipitation due to  
 978 geomagnetic storms: Significance to middle atmosphere ozone chemistry. *Journal*  
 979 *of Geophysical Research*, 115, A11. <https://doi.org/10.1029/2010JA015599><sup>[SEP]</sup>  
 980
- 981 Rozhnoi, A., Solovieva, M.,<sup>[SEP]</sup>Shalimov, S., Ouzounov, D., Gallagher, P., Verth, G.,  
 982 et al. (2020). The effect of the 21 August 2017 total solar eclipse on the phase of  
 983 VLF/LF signals. *Earth and Space Science*, 7, e2019EA000839. <https://doi.org/10.1029/2019EA000839><sup>[SEP]</sup>  
 984  
 985
- 986 Said, R. K., Inan, U. S., & Cummins, K. L. (2010).<sup>[SEP]</sup>Long-range lightning  
 987 geolocation using a VLF radio atmospheric waveform bank. *Journal of*  
 988 *Geophysical Research*, 115, D23. <https://doi.org/10.1029/2010JD013863>  
 989
- 990 Sechrist, C. F. (1968). Interpretation of pre-sunrise electron densities and negative  
 991 ions in the D-region. *Journal of Atmospheric and Terrestrial Physics*, 30(3), 371-  
 992 389. [https://doi.org/10.1016/0021-9169\(68\)90109-8](https://doi.org/10.1016/0021-9169(68)90109-8)  
 993
- 994 Smith, R. A., Coyne, T. N. R., Loch, R. G., & Bourne, I. A. (1967). Small  
 995 perturbation wave interaction in the lower ionosphere, 3, Measurements of electron  
 996 densities, *Ground-based Radio Wave Propagation Studies of the Lower Ionosphere*  
 997 *1 (Compiled by: Belrose, J.S., Bourne, I.A., & Hewitt, L.W.)*, 335-358, Defence  
 998 Research Telecommunications Establishment, Ottawa, Canada.  
 999 [https://publications.gc.ca/collections/collection\\_2019/isde-ised/DR50-2-1967-1-eng.pdf](https://publications.gc.ca/collections/collection_2019/isde-ised/DR50-2-1967-1-eng.pdf)  
 1000
- 1001 Thomas, L. & Bowman, M. R. (1986). A study of pre-sunrise changes in negative  
 1002 ions and electrons in the D-region. *Annales Geophysicae, Series A*, 4(3), 219-227.  
 1003
- 1004 Thomas, L. & Harrison, M. D. (1970). The electron density distributions in the D-  
 1005 region during the night and pre-sunrise period. *Journal of Atmospheric and*  
 1006 *Terrestrial Physics*, 32(1), 1-14. [https://doi.org/10.1016/0021-9169\(70\)90158-3](https://doi.org/10.1016/0021-9169(70)90158-3)  
 1007
- 1008 Thomson, N. R. (1993). Experimental daytime VLF ionospheric parameters. *Journal*  
 1009 *of Atmospheric and Terrestrial Physics*, 55(2), 173–184.  
 1010 [https://doi.org/10.1016/0021-9169\(93\)90122-F](https://doi.org/10.1016/0021-9169(93)90122-F)  
 1011
- 1012 Thomson, N. R. & Clilverd, M. A. (2001), Solar flare induced ionospheric D-region  
 1013 enhancements from VLF amplitude observations. *Journal of Atmospheric and*  
 1014 *Terrestrial Physics*, 63(16), 1729-1737. [https://doi.org/10.1016/S1364-  
 1015 6826\(01\)00048-7](https://doi.org/10.1016/S1364-6826(01)00048-7)  
 1016



- 1017 Thomson, N. R., Rodger, C. J., & Clilverd, M. A. (2005). Large solar flares and  
1018 their ionospheric *D* region enhancements. *Journal of Geophysical Research*,  
1019 *110*(A6), <https://doi.org/10.1029/2005JA011008>  
1020
- 1021 Thomson, N. R., Clilverd, M. A., & McRae, W. M. (2007). Nighttime ionospheric  
1022 *D* region parameters from VLF phase and amplitude. *Journal of Geophysical*  
1023 *Research*, *112*(A7). <https://doi.org/10.1029/2007JA012271>  
1024
- 1025 Thomson, N. R., & McRae, W. M. (2009). Nighttime ionospheric *D* region:  
1026 Equatorial and Non-equatorial, *Journal of Geophysical Research*, *114*(A8).  
1027 <https://doi.org/10.1029/2008JA014001>  
1028
- 1029 Thomson, N. R. (2010). Daytime tropical *D* region parameters from short path VLF  
1030 phase and amplitude. *Journal of Geophysical Research*, *115*(A9).  
1031 <https://doi.org/10.1029/2010JA015355>  
1032
- 1033 Thomson, N. R., Clilverd, M. A., & Rodger, C. J. (2011a). Daytime midlatitude  
1034 *D* region parameters at solar minimum from short-path VLF phase and amplitude,  
1035 *Journal of Geophysical Research*, *116*(A3). <https://doi.org/10.1029/2010JA016248>  
1036
- 1037 Thomson, N. R., Rodger, C. J., & Clilverd, M. A. (2011b). Daytime *D* region  
1038 parameters from long-path VLF phase and amplitude, *Journal of Geophysical*  
1039 *Research*, *116*(A11). <https://doi.org/10.1029/2011JA016910>  
1040
- 1041 Thomson, N. R., Rodger, C. J., & Clilverd, M. A. (2012). Tropical daytime lower  
1042 *D* region dependence on sunspot number. *Journal of Geophysical Research*  
1043 *117*(A10). <https://doi.org/10.1029/2012JA018077>  
1044
- 1045 Thomson, N. R., Clilverd, M. A., & Rodger, C. J. (2014). Low-latitude ionospheric  
1046 *D* region dependence on solar zenith angle. *Journal of Geophysical Research:*  
1047 *Space Physics* *119*(8). <https://doi.org/10.1002/2014JA020299>  
1048
- 1049 Thomson, N. R., Clilverd, M. A., & Rodger, C. J. (2017). Midlatitude ionospheric  
1050 *D* region: Height, sharpness, and solar zenith angle. *Journal of Geophysical*  
1051 *Research: Space Physics*, *122*(8), 8933-8946.  
1052 <https://doi.org/10.1002/2017JA024455>  
1053
- 1054 Thomson, N. R., Clilverd, M. A., & Rodger, C. J. (2018). Quiet daytime Arctic  
1055 ionospheric *D* region. *Journal of Geophysical Research: Space Physics*, *123*,  
1056 9726–9742. <https://doi.org/10.1029/2018JA025669>  
1057
- 1058 Thomson, N. R., Clilverd, M. A., Brundell, J. B., & Rodger, C. J. (2021). Quiet night  
1059 Arctic ionospheric *D* region characteristics. *Journal of Geophysical Research:*  
1060 *Space Physics*, *126*, e2020JA029043. <https://doi.org/10.1029/2020JA029043>  
1061
- 1062 Thomson, N. R., Clilverd, M. A., & Rodger, C. J. (2022). Ionospheric *D* region: VLF-  
1063 measured electron densities compared with rocket-based FIRI-2018 model.  
1064 *Journal of Geophysical Research: Space Physics*, *127*, e2022JA030977.  
1065 <https://doi.org/10.1029/2022JA030977>  
1066

- 1067 Turco, R. P. (1974). A discussion of possible negative ion detachment mechanisms in  
1068 the sunrise D region. *Radio Science* 9(7), 655-658.  
1069 <https://doi.org/10.1029/RS009i007p00655>  
1070
- 1071 Verronen, P. T., Ulich, Th., Turunen, E., & Rodger, C. J. (2006). Sunset transition of  
1072 negative charge in the D-region ionosphere during high-ionization conditions.  
1073 *Annales Geophysicae*, 24, 187–202. <https://doi.org/10.5194/angeo-24-187-2006>  
1074
- 1075 Wait, J. R., & Spies, K. P. (1964). Characteristics of the Earth-ionosphere waveguide  
1076 for VLF radio waves. *NBS Tech. Note 300*, Natl. Bur. of Stand., Boulder, Colo.  
1077 [https://www.govinfo.gov/app/details/GOVPUB-C13-  
1078 1fc83a916d87542f34917847f89b9f0b](https://www.govinfo.gov/app/details/GOVPUB-C13-1fc83a916d87542f34917847f89b9f0b)  
1079
- 1080 Worthington, E. R., & Cohen, M. B. (2021). The estimation of D-region electron  
1081 densities from trans-ionospheric very low frequency signals. *Journal of*  
1082 *Geophysical Research: Space Physics*, 126, e2021JA029256.  
1083 <https://doi.org/10.1029/2021JA029256>  
1084
- 1085 Xu, W., Marshall, R. A., Bortnik, J., & Bonnell, J. W. (2021). An electron density  
1086 model of the D- and E-region ionosphere for transionospheric VLF propagation.  
1087 *Journal of Geophysical Research: Space Physics*, 126, e2021JA029288.  
1088 <https://doi.org/10.1029/2021JA029288>  
1089
- 1090 Zhou, X., Wang, J., Ma, D., Huang, Q., & Xiao, F. (2021). A method for determining  
1091 D region ionosphere reflection height from lightning skywaves. *Journal of*  
1092 *Atmospheric and Solar-Terrestrial Physics*, 221, 105692.  
1093 <https://doi.org/10.1016/j.jastp.2021.105692>  
1094

1095 **Figure Captions**

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1097 **Figure 1.** VLF paths used here to determine ionospheric *D* region characteristics

1098 particularly around dawn and dusk.

1099 **Figure 2.** (Top two panels) Phase (degrees) and amplitude (dB) observed at

1100 Oakington, Cambridge, UK, versus time (0-12 UT) received from the 45.9 kHz

1101 transmitter NSY near Niscemi, Sicily, 30 July – 4 August 2005.

1102 (Middle two panels) Same data as in the top two panels but for an expanded time

1103 scale about dawn (3.2-4.6 UT)

1104 (Bottom two panels) ModeFinder calculations for this path of the phase (degrees) and

1105 amplitude (dB > 1  $\mu$ V/m) as functions of height,  $H'$  and sharpness,  $\beta$ .

1106 **Figure 3.** Height parameter,  $H'$  versus SZA inferred from VLF observations.

1107 (Top two panels) Mid-latitude values of  $H'$  versus SZA for morning/dawn (left) and

1108 afternoon/dusk (right).

1109 (Middle two panels) Low latitude values of  $H'$  versus SZA for morning/dawn (left)

1110 and afternoon/dusk (right).

1111 (Bottom two panels) High mid-latitude values of  $H'$  versus SZA for morning/dawn

1112 (left) and afternoon/dusk (right).

1113 **Figure 4.** Sharpness parameter,  $\beta$ , versus SZA inferred from VLF observations.

1114 (Top two panels) Mid-latitude values of  $\beta$  versus SZA for morning/dawn (left) and

1115 afternoon/dusk (right).

1116 (Middle two panels) Low latitude values of  $\beta$  versus SZA for morning/dawn (left) and

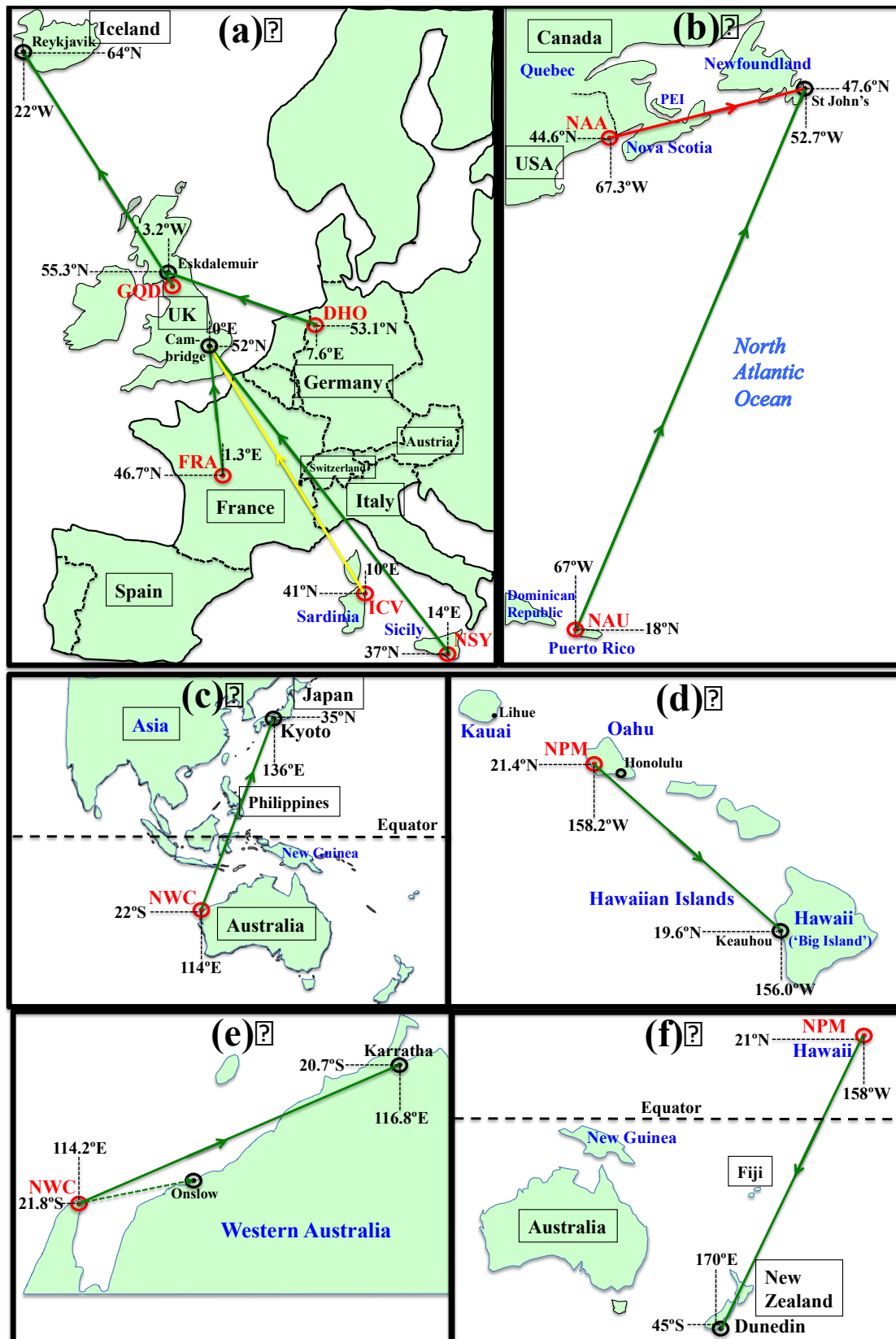
1117 afternoon/dusk (right).

1118 (Bottom two panels) High mid-latitude values of  $\beta$  versus SZA for morning/dawn

1119 (left) and afternoon/dusk (right).



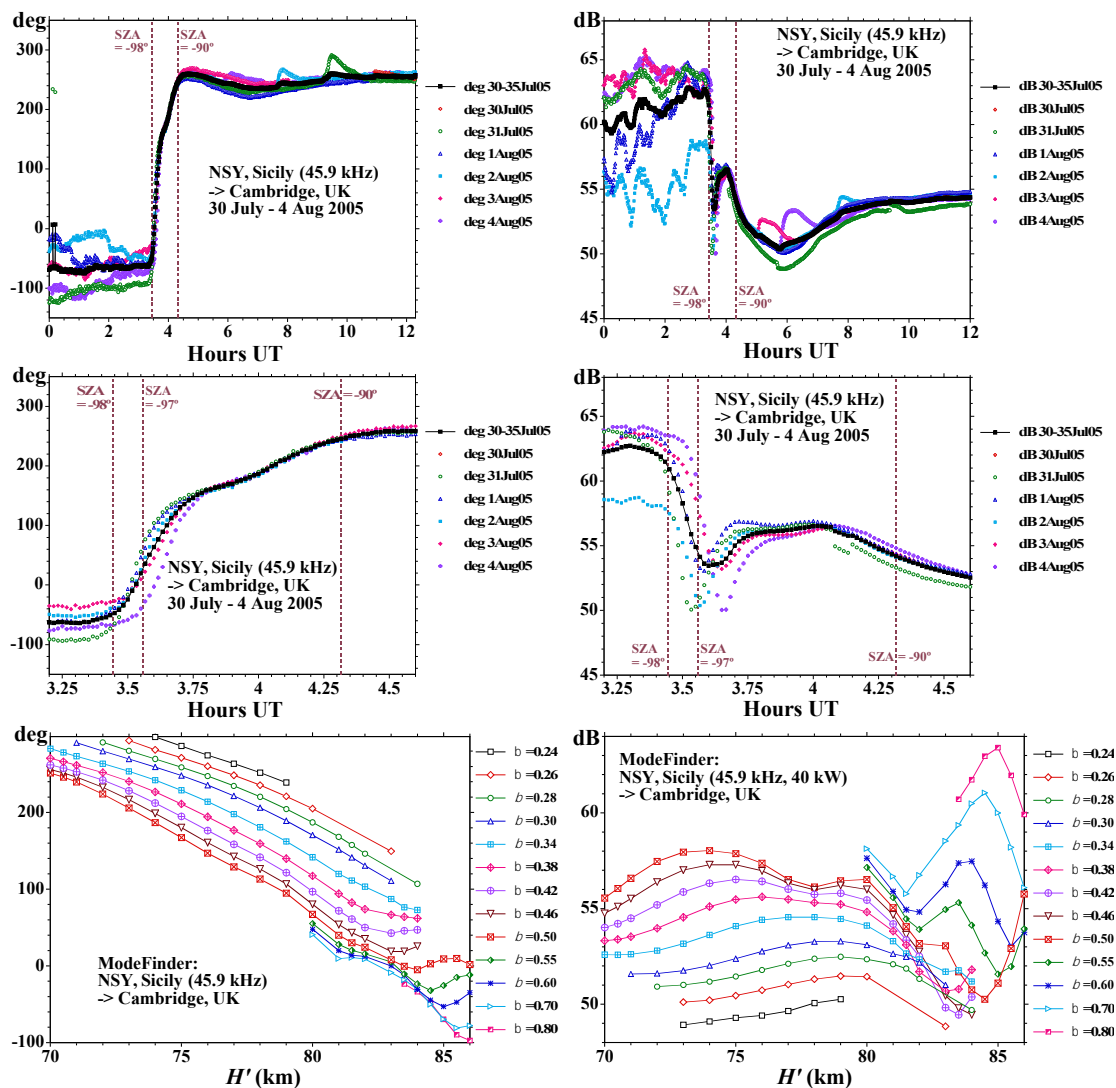
1120 **Figure 5.** Electron number density height profiles at dawn ( $SZA = -90^\circ$ ) at mid-  
1121 latitude. The green circles are from the data, reported by Mechtly & Smith (1970),  
1122 from a rocket flown at Wallops Island ( $38^\circ N$ ), Virginia, USA, on 24 July 1968. The  
1123 blue circles are the same rocket data but raised in height by 2 km (see text). The  
1124 heavy black line with square plot symbols is from the NSY-Cambridge VLF path,  
1125 30 July – 4 Aug 2005, where  $H' = 72.75$  km and  $\beta = 0.378$  km $^{-1}$  matched the data (at  
1126  $SZA = -90^\circ$ ). The orange line at 71 km is indicating the greatest height at which the  
1127 VLF propagation on this path is sensitive to the electron number density. The purple  
1128 diamonds are from the FIRI-2018 rocket-based model at dawn ( $SZA = -90^\circ$ ), for July at  
1129 latitude  $45^\circ$ .  
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1133 **Figure 1.** VLF paths used here to determine ionospheric *D* region characteristics  
1134 **particularly** around dawn and dusk.

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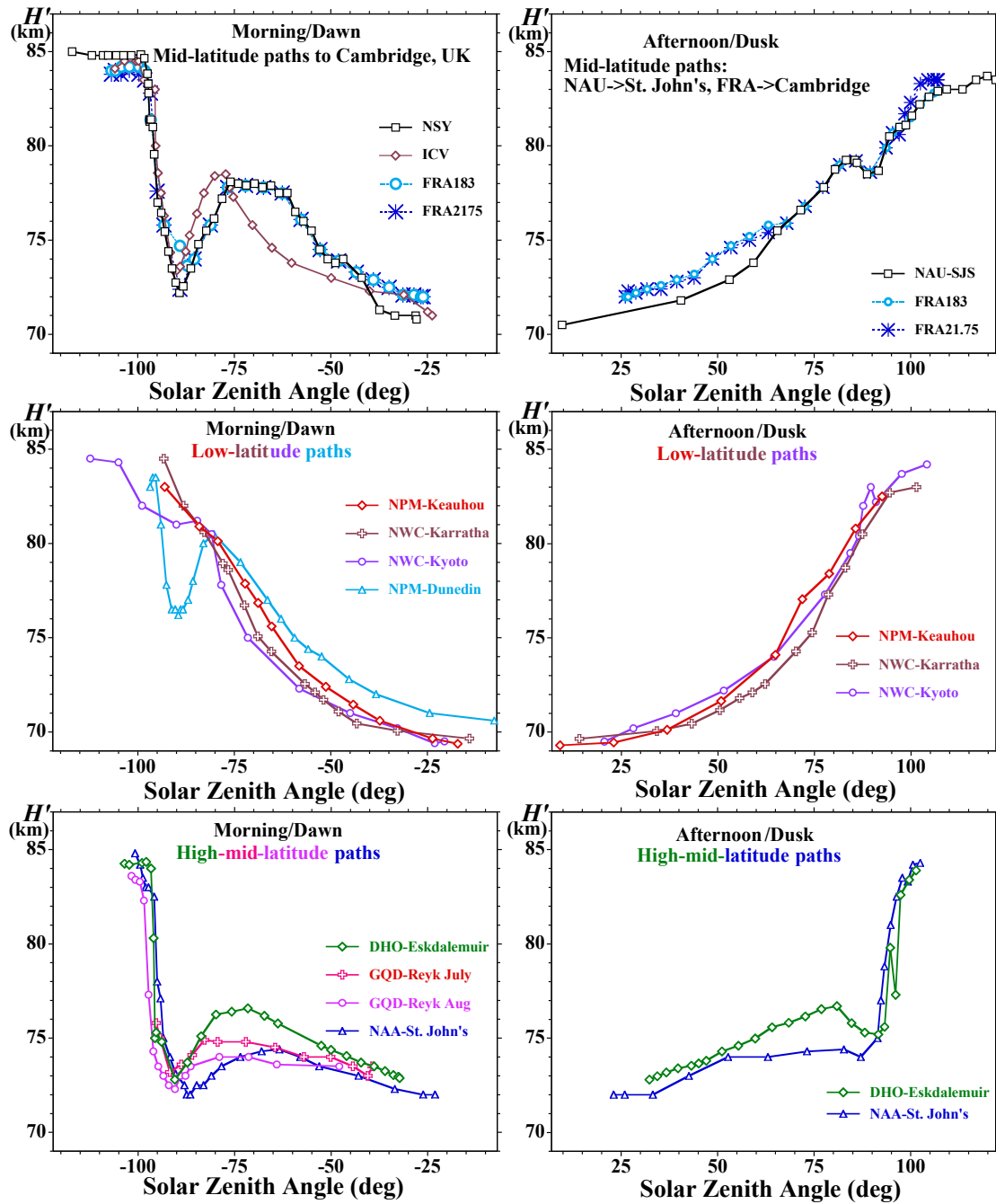
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1142 scale about dawn (3.2-4.6 UT)

1143 (Bottom two panels) ModeFinder calculations for this path of the phase (degrees) and

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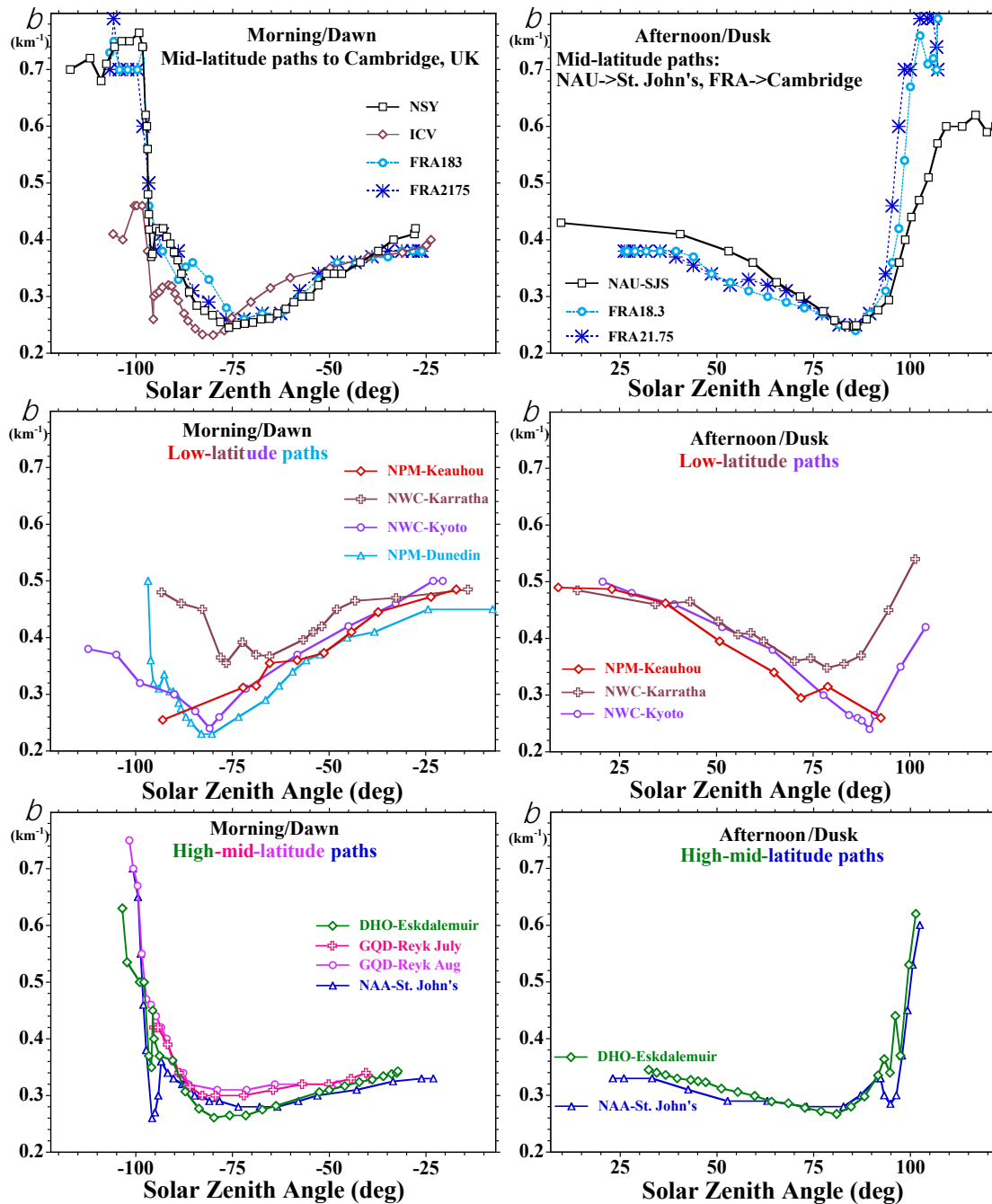
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1154 (left) and afternoon/dusk (right).



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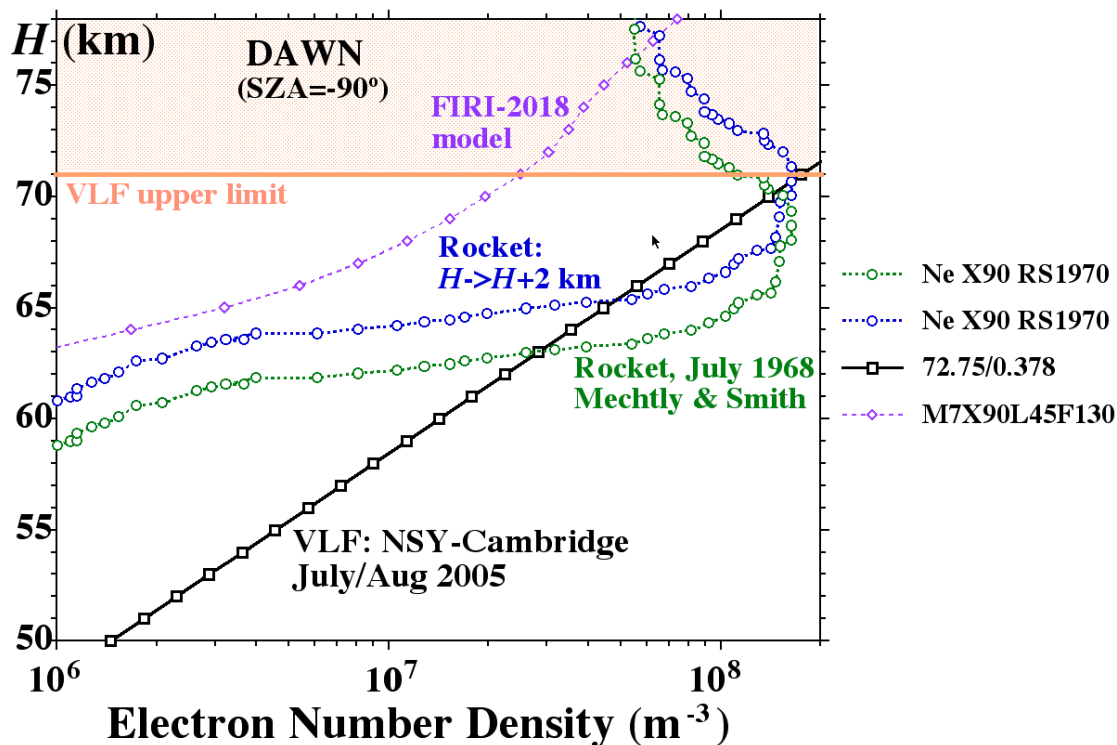
1160 (Middle two panels) Low latitude values of  $\beta$  versus SZA for morning/dawn (left) and

1161 afternoon/dusk (right).

1162 (Bottom two panels) High mid-latitude values of  $\beta$  versus SZA for morning/dawn

1163 (left) and afternoon/dusk (right).

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**Figure 1.**

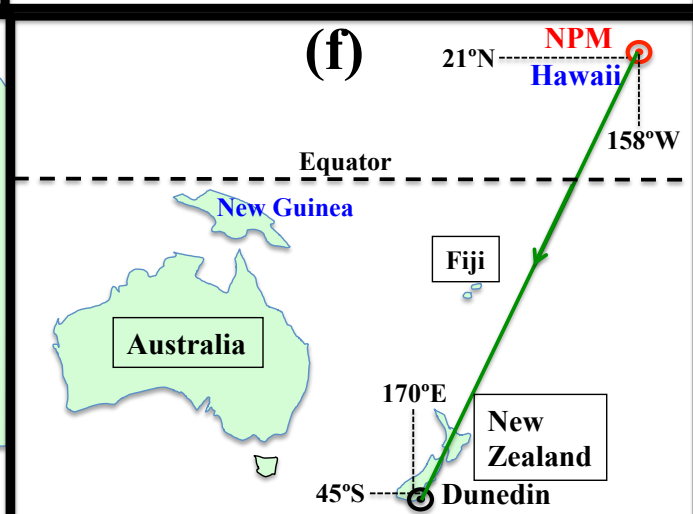
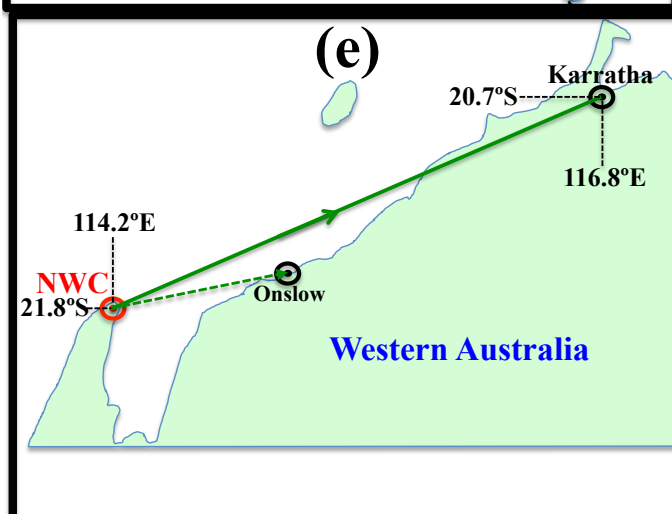
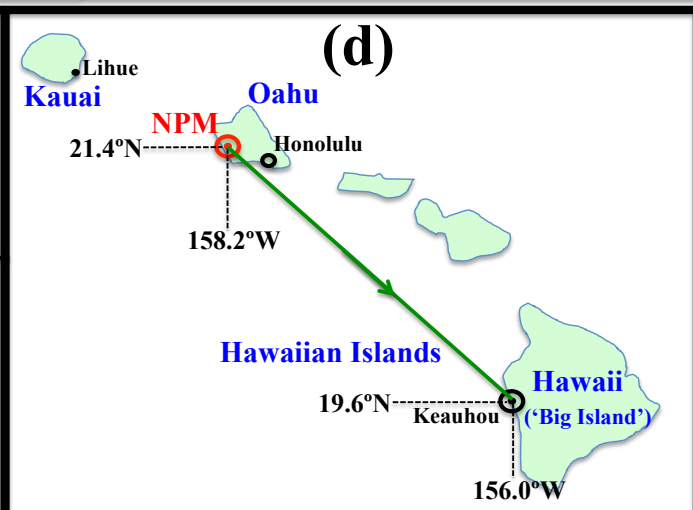
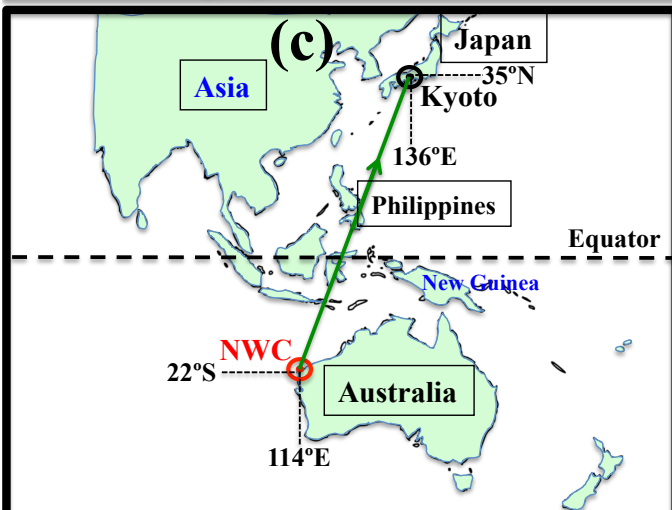
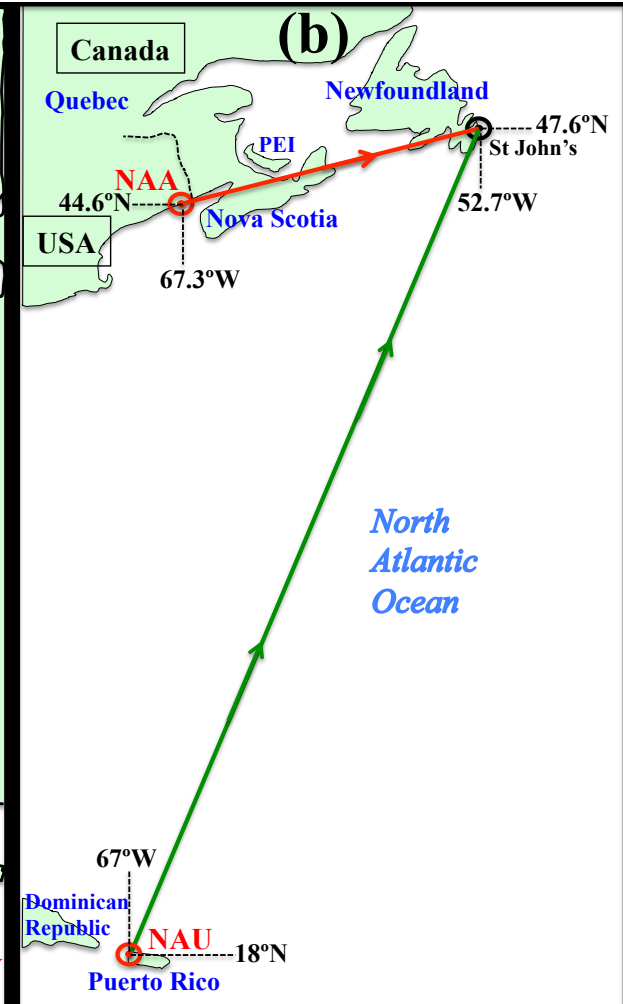
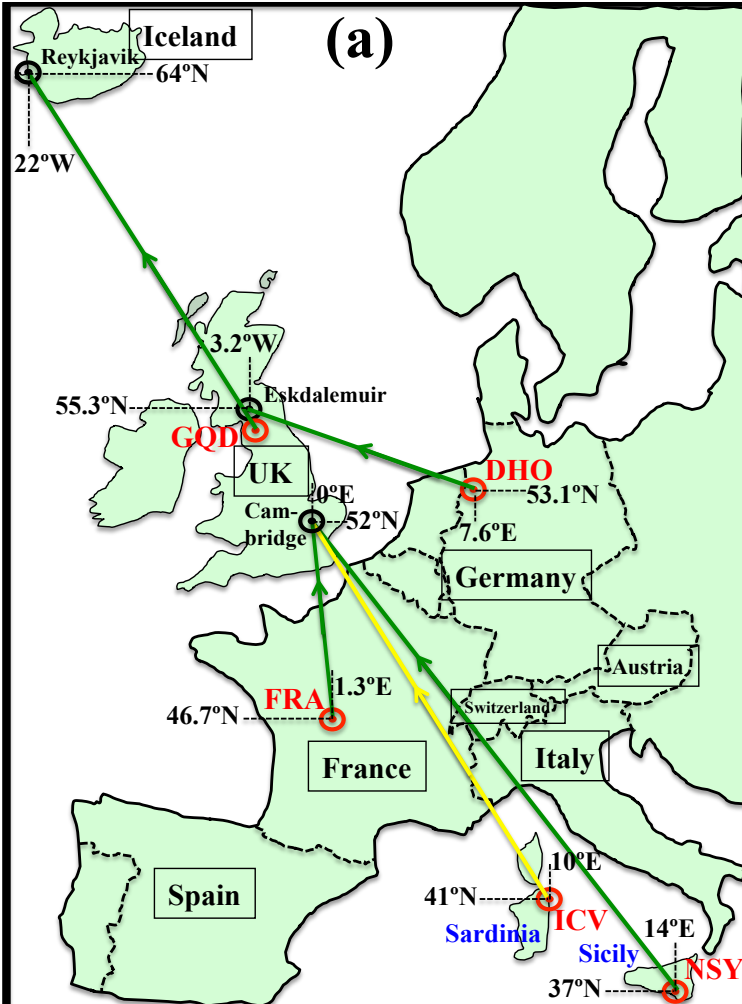




Figure 2.

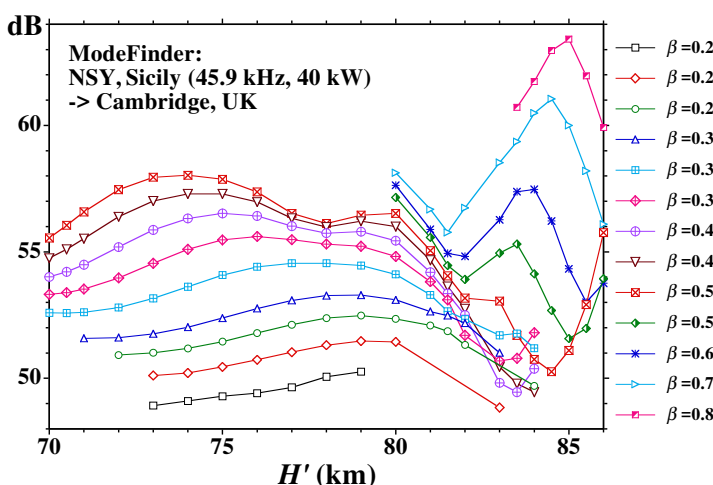
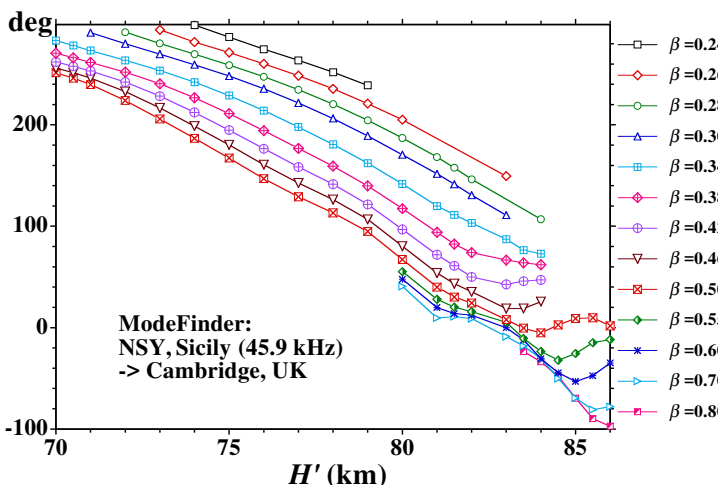
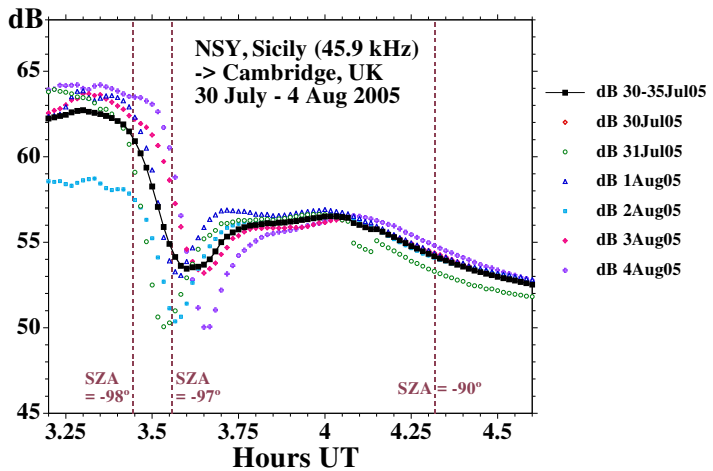
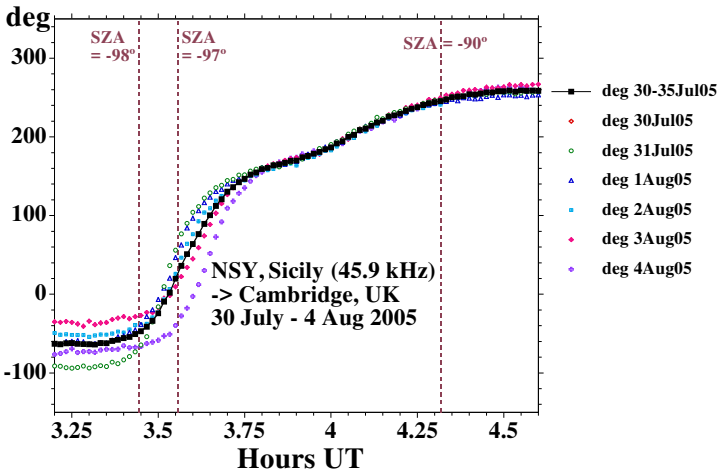
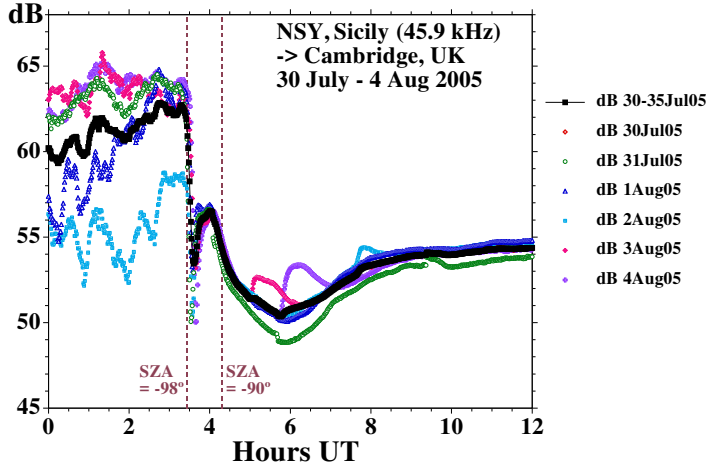
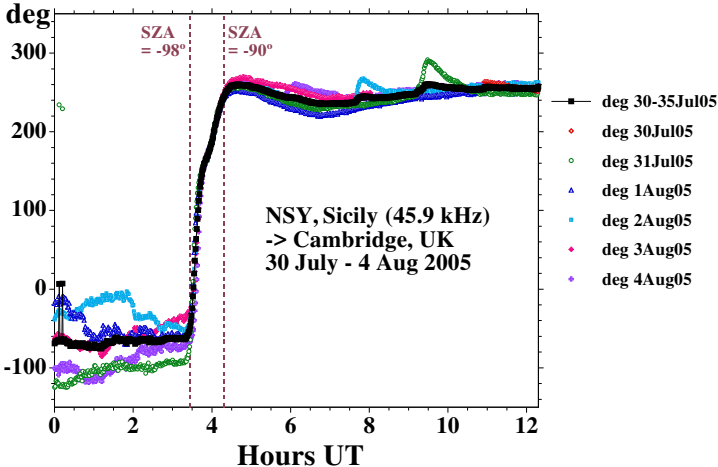


Figure 3.

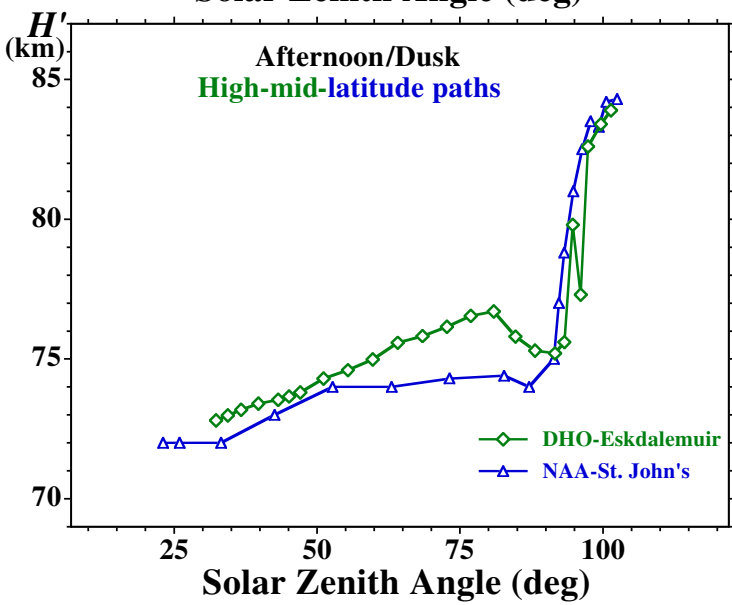
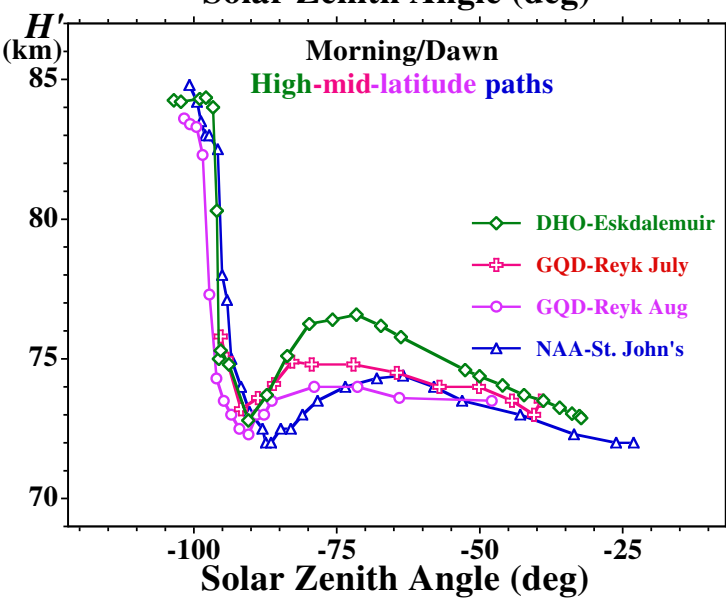
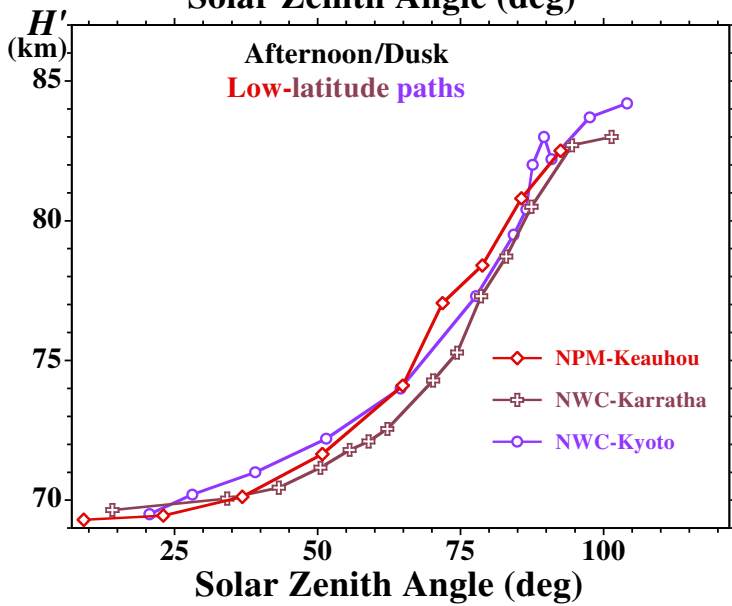
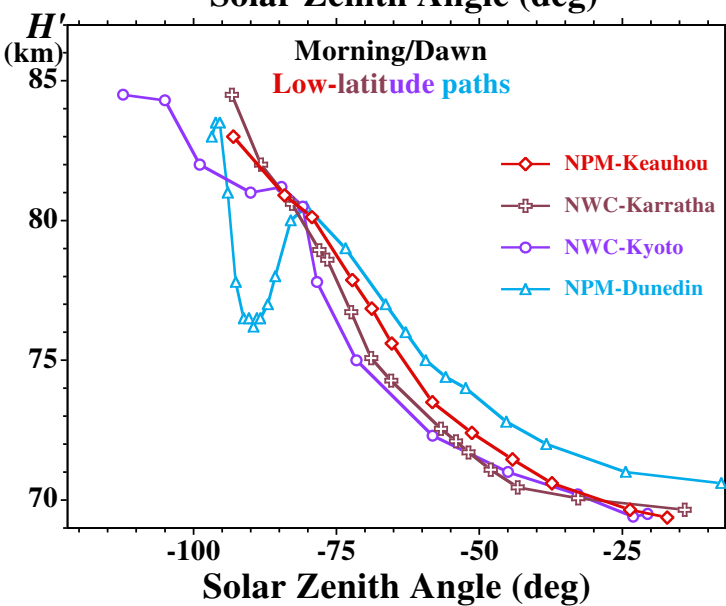
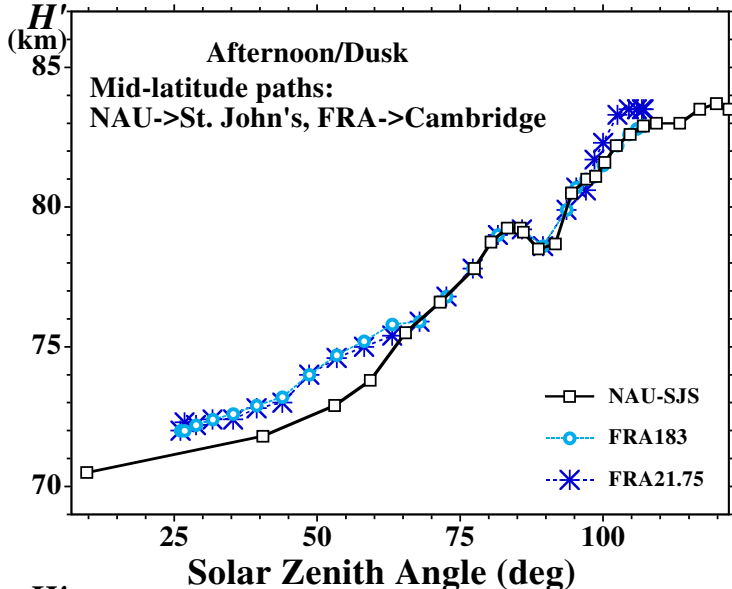
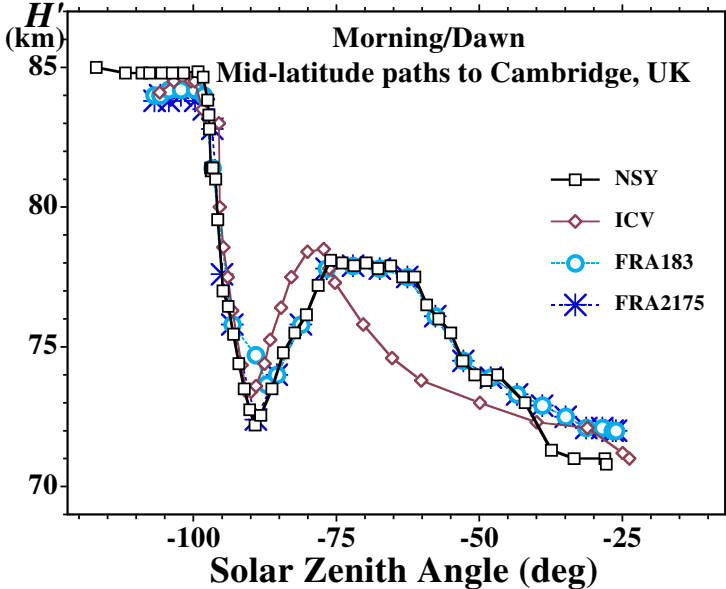


Figure 4.

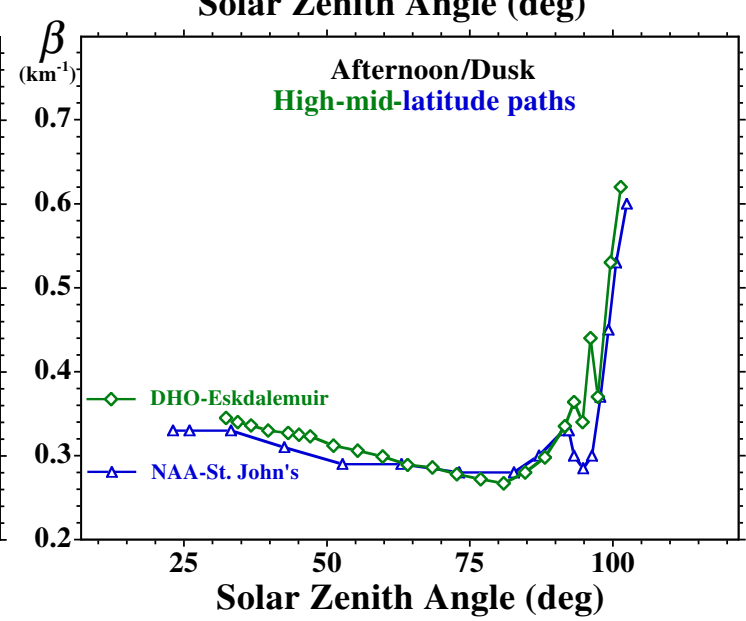
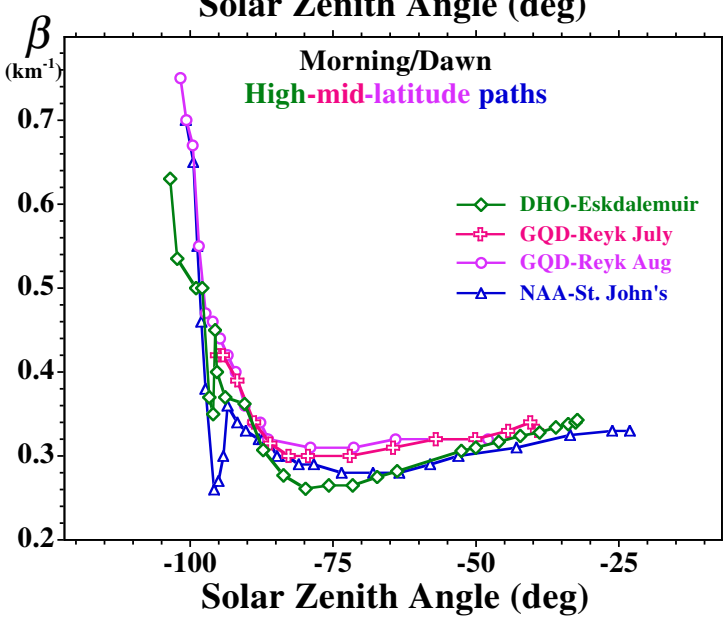
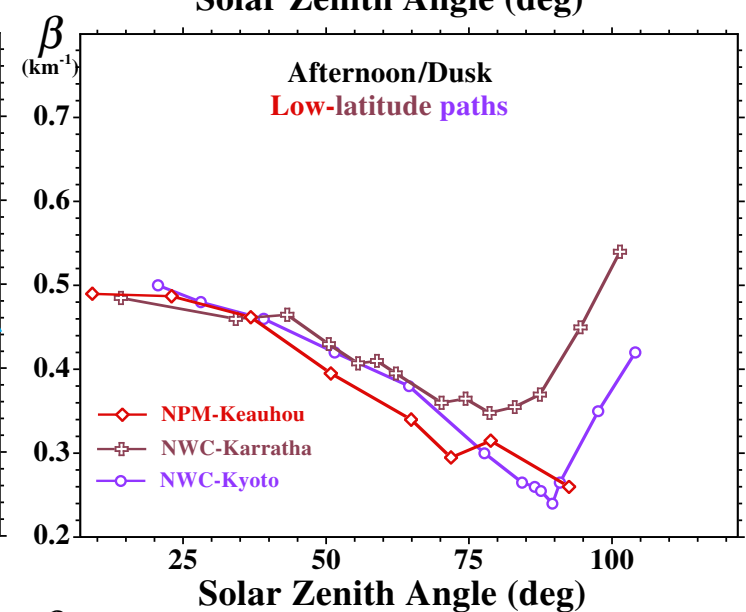
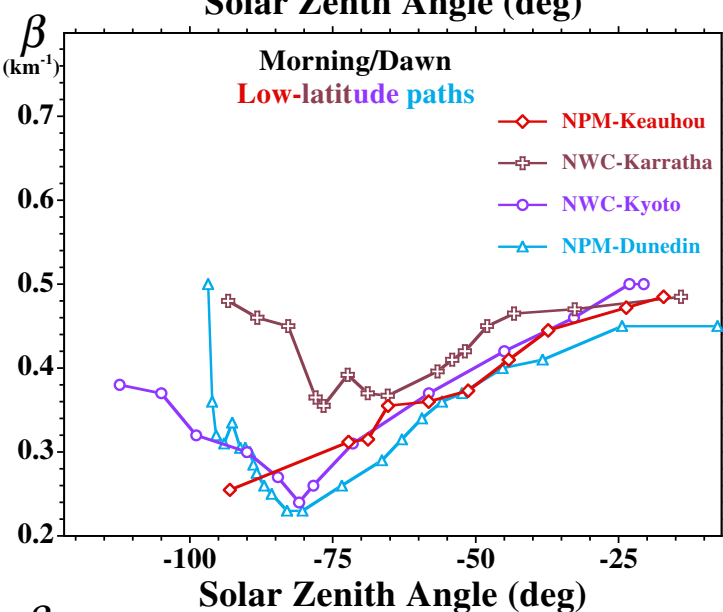
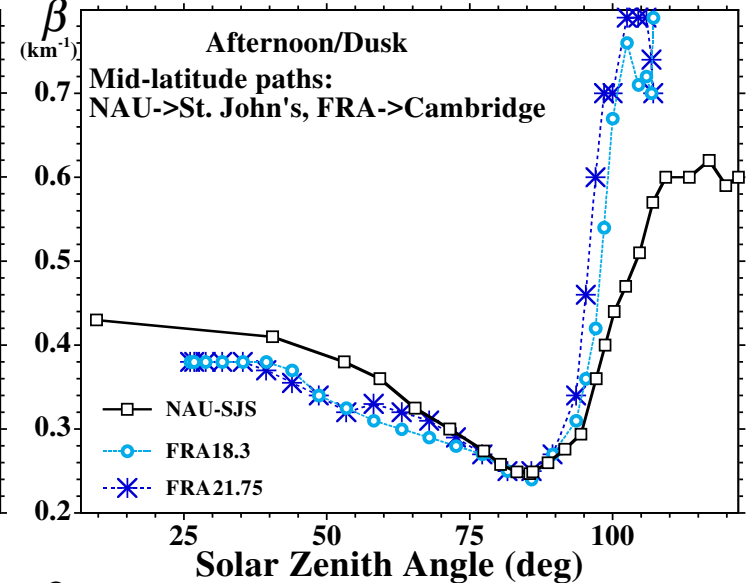
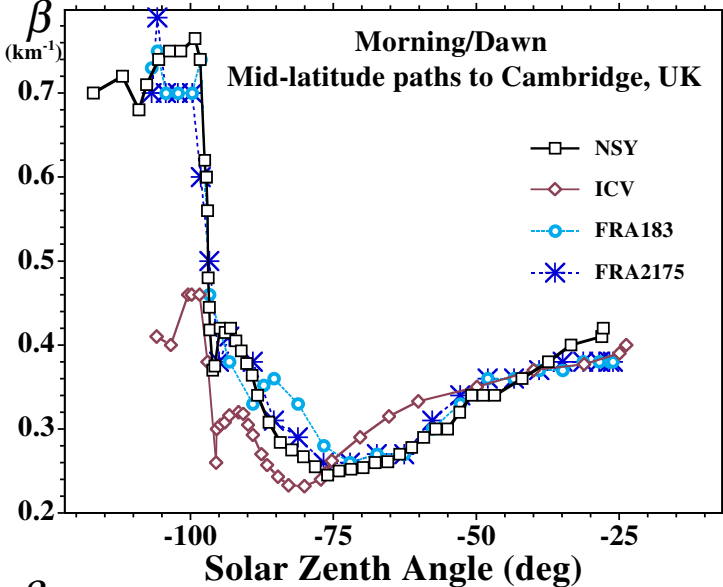


Figure 5.

